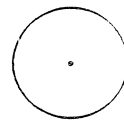
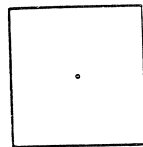
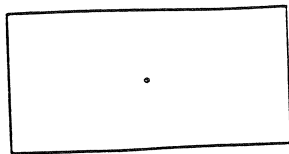
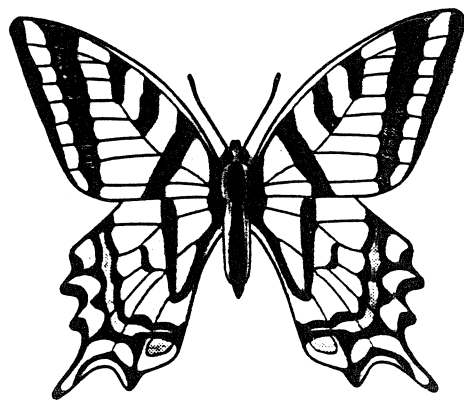


Symmetries in quantum mechanics

What is a symmetry?



A **symmetry** (or symmetry transformation) of an object is a transformation which maps the object into itself. If an object admits a certain symmetry it is said to be **invariant** under the transformation.

The set of all symmetry transformations of an object represents a group: the **symmetry group** of the object.

Symmetries of the laws of Nature

A **symmetry transformation** of a physical law is a change of the variables and/or the space-time coordinates (in terms of which it is formulated) such that the equations describing the law have the same form in terms of the new variables and coordinates as they had in terms of the old one: the equations are said to be **covariant** w.r.t. the symmetry transformation.

Ex. the wave equation

$$\left\{ \square \psi(\vec{r}, t) = 0 \right.$$

\uparrow $\square = -\nabla^2 + \frac{1}{c^2} \frac{\partial^2}{\partial t^2}$

Lorentz boost

$$x' = \gamma(x - vt), \quad y' = y, \quad z' = z, \quad t' = \gamma\left(t - \frac{v}{c^2}x\right)$$

$$\gamma = (1 - (v/c)^2)^{-1/2}$$

\Downarrow Lorentz

$$\square' = \square, \quad \psi'(\vec{r}', t') = \psi(\vec{r}, t)$$

\Rightarrow THE WAVE EQUATION IS COVARIANT UNDER LORENTZ BOOSTS

Space-time vs. internal symmetries



SYMMETRIES

SPACE-TIME
SYMMETRIES

IN INVARIANCE UNDER

TRANSLATIONS

ROTATIONS

⋮

LORENZ BOOSTS

⋮

SCALE CHANGES

⋮

INTERNAL
SYMMETRIES

IN INVARIANCE UNDER

GLOBAL PHASE TRANSFORMATIONS

GAUGE TRANSFORMATIONS

⋮

Space-time symmetries

translation in time : $t \rightarrow t + a_0$

translations in space : $\vec{r} \rightarrow \vec{r} + \vec{a}$

rotations in space : $\vec{r} \rightarrow \mathbb{R} \vec{r}$

**Induced transformations
in Hilbert space**

Galilei transformations : $\vec{r} \rightarrow \vec{r} - \vec{v}t$ (non-relativistic)

Lorentz boosts : previous page (relativistic)

Lorentz group \mathcal{L} = Lorentz boosts
+ rotations in space
+ parity + time reversal

$$\mathcal{P} = \mathcal{L} \otimes \mathcal{T}_4 \quad \text{POINCARÉ' GROUP}$$

EXTENSIONS

CONFORMAL
GROUP

SUPER POINCARÉ'

.....

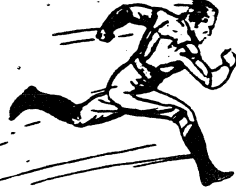
$$\mathcal{P} \otimes \{\text{SCALE CHANGES}\}$$

$$\begin{cases} x \rightarrow \lambda x \end{cases}$$

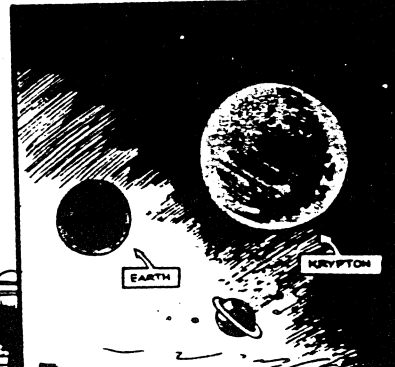
$$\mathcal{P} \otimes \{\text{SUSY}\}$$

↳
BOSONS ↔ FERMIONS

SCIENTIFIC EXPLANATION OF **SUPERMAN'S AMAZING STRENGTH --!**



SUPERMAN CAME TO EARTH FROM THE PLANET **KRYPTON**, WHOSE INHABITANTS HAD EVOLVED, AFTER MILLIONS OF YEARS, TO PHYSICAL PERFECTION!



THE SMALLER SIZE OF OUR PLANET, WITH ITS SLIGHTER GRAVITY PULL, ASSISTS **SUPERMAN'S** TREMENDOUS MUSCLES IN THE PERFORMANCE OF MIRACULOUS FEATS OF STRENGTH!

EVEN UPON OUR WORLD TODAY EXIST CREATURES POSSESSING **SUPER-STRENGTH!**



THE LOWLY ANT CAN SUPPORT WEIGHTS HUNDREDS OF TIMES ITS OWN.



THE GRASSHOPPER LEAPS WHAT TO MAN WOULD BE THE SPACE OF SEVERAL CITY BLOCKS!

IT IS NOT TOO FAR FETCHED TO PREDICT THAT SOME DAY OUR VERY OWN PLANET MAY BE PEOPLED ENTIRELY BY **SUPERMEN**!



IS THIS RIGHT ?

DO WE HAVE SCALE INVARIANCE IN A GRAVITATIONAL ENVIRONMENT ?

ARE SYMMETRIES ALWAYS MANIFEST ?

Internal symmetries

Global phase symmetries

$$H\psi = i\hbar \frac{\partial \psi}{\partial t} \quad \text{INVARIANT UNDER } \psi \rightarrow e^{i\alpha} \psi, \quad \alpha \in \mathbb{R}$$

Gauge symmetries (local phase symmetries)

$$H\psi = i\hbar \frac{\partial \psi}{\partial t} \quad \text{INVARIANT UNDER } \psi \rightarrow e^{i\alpha(\vec{r}, t)} \psi$$

↑
local phase

PROVIDED ONE PROPERLY
INTRODUCES A GAUGE FIELD

$$A_\mu$$

IN H!

MAXWELL THEORY
("U(1) GAUGE SYMMETRY")

$$: A_\mu = (\phi, -\vec{A})$$

$$\vec{E} = -\nabla\phi - \frac{\partial \vec{A}}{\partial t}$$

$$\vec{B} = \nabla \times \vec{A}$$

WHY ARE SYMMETRIES USEFUL IN PHYSICS ?

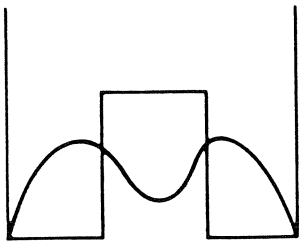
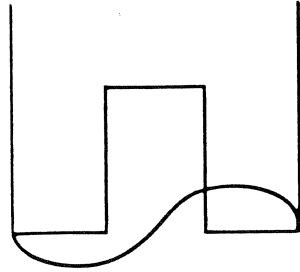
NOETHER'S THEOREM

Covariance of the equations of motion with respect to a continuous transformation with n parameters implies the existence of n conserved quantities ('conserved charges' or 'integrals of motion'), i.e. it implies conservation laws.

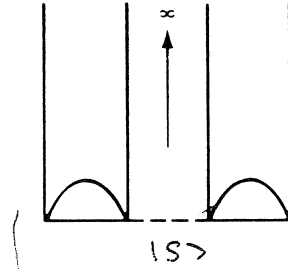
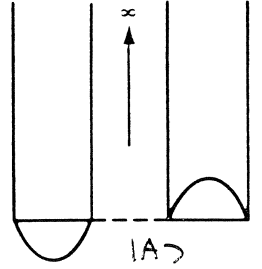
Property	Invariance of equations	Conserved quantity
homogeneity of time	time translation invariance	energy
homogeneity of space	translational invariance	momentum
isotropy of space	rotational invariance	angular momentum

DISCRETE SYMMETRIES ARE ALSO USEFUL ...

SYMMETRIC DOUBLE-WELL POTENTIAL

Symmetric $|S\rangle$ Anti-symmetric $|A\rangle$

$$E_A > E_S$$

 $|S\rangle$  $|A\rangle$

$$E_A = E_S$$

$$\begin{aligned}
 |R\rangle &= \frac{1}{\sqrt{2}} (|S\rangle + |A\rangle) \longrightarrow \frac{1}{\sqrt{2}} (e^{-iE_S t/\hbar} |S\rangle + e^{-iE_A t/\hbar} |A\rangle) \\
 &= \frac{1}{\sqrt{2}} e^{-iE_S t/\hbar} (|S\rangle + e^{-i(E_A - E_S)t/\hbar} |A\rangle)
 \end{aligned}$$

$$|L\rangle = \frac{1}{\sqrt{2}} (|S\rangle - |A\rangle) \longrightarrow \frac{1}{\sqrt{2}} e^{-iE_S t/\hbar} (|S\rangle - e^{-i(E_A - E_S)t/\hbar} |A\rangle)$$

$V(x) \rightarrow \infty \Rightarrow$ NO TUNNELING

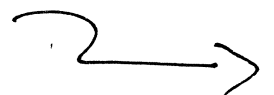
$|S\rangle$ AND $|A\rangle$ DEGENERATE \Rightarrow $|R\rangle$ AND $|L\rangle$

ALSO ENERGY EIGENSTATES

THE GROUND STATE MAY BREAK THE PARITY SYMMETRY
ALTHOUGH THE HAMILTONIAN IS INVARIANT UNDER
PARITY

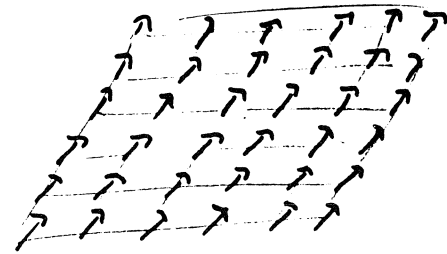
EXAMPLE OF SPONTANEOUSLY BROKEN SYMMETRY

... THE GROUND STATE HAS LESS SYMMETRY
THAN THE HAMILTONIAN



Paradigm of a **broken symmetry**: **FERROMAGNET**

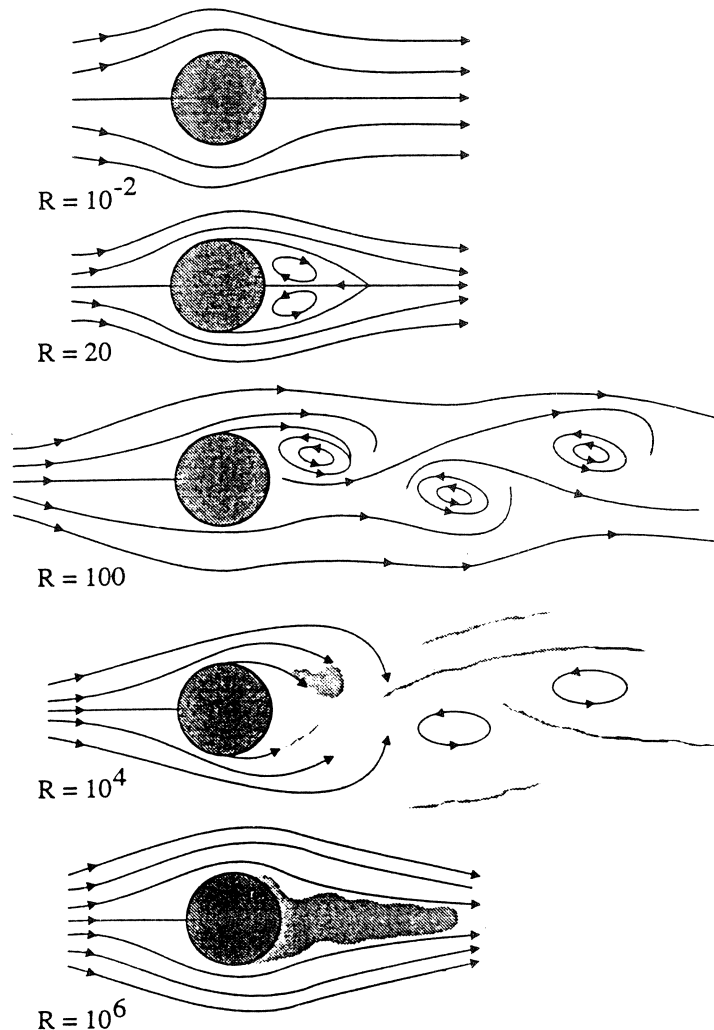
Hamiltonian $H = -J \sum_{\langle ij \rangle} \vec{S}_i \cdot \vec{S}_j$



Ground state $|\psi\rangle = |\uparrow\uparrow\uparrow \dots \uparrow\rangle$

breaks SU(2) ----> low-lying excitations are *Goldstone bosons*

Another example (classical physics!) is **TURBULENCE**, where different symmetries are broken spontaneously at different scales

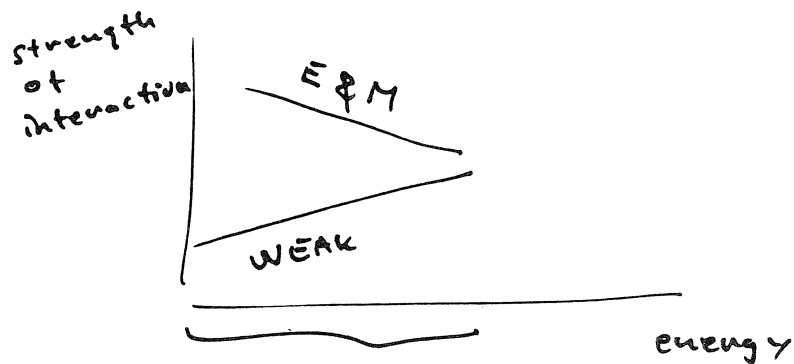


Yet another example of a broken symmetry...

Electroweak theory (Glashow, Weinberg, Salam, 1960s)

$U(1) \times SU(2)$ gauge symmetry spontaneously broken

W^\pm, Z^0 acquire masses ("Anderson-Higgs mechanism", like a photon in a superconductor!)



$$\langle 0 | \psi | 0 \rangle \neq 0$$

↑ "Higgs field"

$$Cf \quad \langle \psi | \vec{S} | \psi \rangle \neq 0$$

in a ferromagnet

THE DIFFERENT PHASES OF MATTER CAN BE CLASSIFIED ACCORDING TO WHICH SYMMETRIES THEY "SPONTANEOUSLY" BREAK

EXAMPLE :

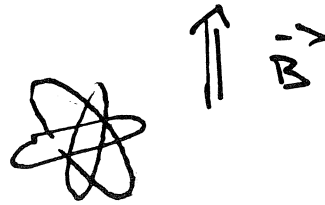
TABLE OF BROKEN-SYMMETRY PHENOMENA

Phenomenon	High Phase	Low Phase	Order Parameter (constant of motion)	Order-Parameter Dimensionality $T \rightarrow 0$	T_c	Common Transition Type (Can Always Be First Order)	"Goldstone Bosons" (or "Higgs" Bosons)	Fluctuations	Collective Hydrodynamic Modes	Generalized Rigidity Phenomenon	Long-Range Forces Due to General	Rigidity	Singularities
Ferroelectricity (Pyroelectricity)	Non-Polar crystal	Polar crystal	\vec{P} (no)	1	1 or 3	2nd or 1st nearly 2nd	no (optical phonons)	Soft Modes	No	Ferroelectric hysteresis	No		Domain walls (thin)
Ferromagnetism	Paramagnet	Ferromagnet	\vec{M} (yes)	1, often ≈ 3	1 or 3	2nd (time reversal)	Spin waves, one branch $\omega \propto k^2 (+ \text{const})$	Spin waves	No ?	Permanent magnetism: hysteresis	Suhl-Nakamura		Domain walls (very mobile, $\approx 1 \text{ cm}^2$)
Antiferromagnetism	Paramagnet	Antiferromagnet	$\vec{M}_{\text{sublattice}}$ (no)	1, often ≈ 3	1 or 3	2nd (or first)	Spin waves, 2 branches $\omega \propto k + \text{const}$ ("Fermions" in metal case?)	Spin waves (diffusion?)	No?	subtle effects in A.F. resonance	Suhl-Nakamura		Domain walls
Superconductivity	Normal Metal	Superconductor	$\langle \psi^* \psi \rangle$ $= \text{Fe}^{1/2}$ (no)	2	2	2nd (Gauge; no 3rd order terms)	no (plasmons)	diffusive fluctuations of gap	Mostly not	superconductivity	No: Penetration depth		Flux lines (or normal domains in Type I)
He II	Normal liquid	Superfluid	$\langle \psi \rangle$ (no)	2	2	2nd (Gauge; no 3rd order terms)	phonons (1 branch)	diffusive fluctuations of $\langle \psi \rangle$	2nd sound, especially	superfluidity	Yes: vortex lines unscreened		vortex lines
Nematic liquid crystal	Normal liquid	oriented liquid	\vec{d} (directrix) (no)	3	3	2nd	not stable at $T = 0$ (yes in principle)	Yes	usually overdamped	various peculiar properties: orientation elasticity	Yes		disclinations, points
Cholesteric, Smectic liquid crystal	Nematic liquid	Density wave	$\rho(Q)$ (no)	> 3	> 3	2nd or 1st		Yes			Yes		disclinations, points and dislocations
Crystal	liquid	solid	$\rho_{\vec{r}}$, all G on reciprocal lattice (no)	3 (2 orient., 1 phase) at least	3	1st	yes: 3 kinds 2 transverse 1 longitudinal	phonons	2nd sound in solid, etc.	rigidity	Yes: elasticity effects		dislocations, grain boundaries, points (vacancies, interstitials)
He 3	normal liquid	anisotropic superfluid	$d_{ij} = \langle \psi \psi \rangle_{M_1, M_2}$ (no)	at least 3	18	2nd	yes, several kinds	yes, complex	probably	superfluidity and orientation elasticities	Yes		Vortex lines, dislocations, point defects
CDW's	normal electron gas (2-dim)	Incommensurate Density wave + commensurate	$\rho_{\vec{r}}$, G on triangular superlattice (no)	2	2	1st	Yes, phasons		Yes	Yes, NbSe ₂ sliding CDW's			discommensurations, dislocations

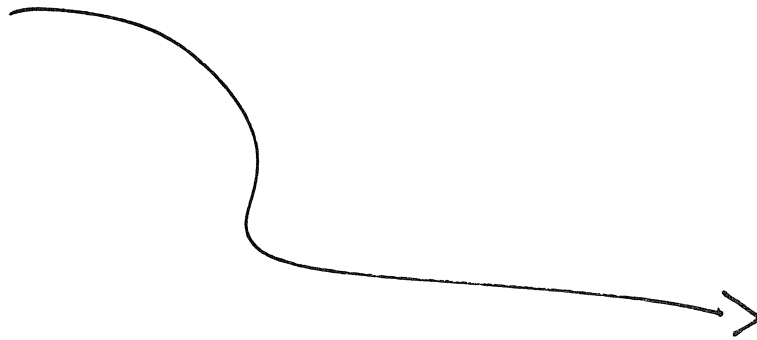
Symmetries can also be **explicitly broken**

examples:

atom in a magnetic field



parity!

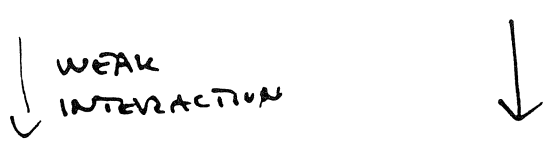


CAN PARITY BE EXPLICITLY BROKEN ?

$$[\pi, H] \neq 0 \quad ?$$

YES! IN WEAK INTERACTIONS!

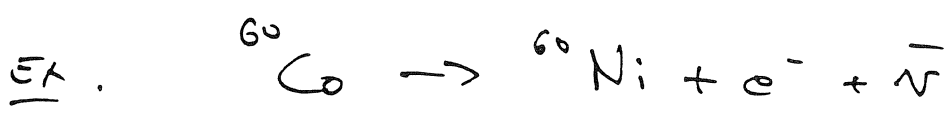
$$|\psi(0)\rangle \xleftrightarrow{\text{PARITY}} \pi |\psi(0)\rangle$$



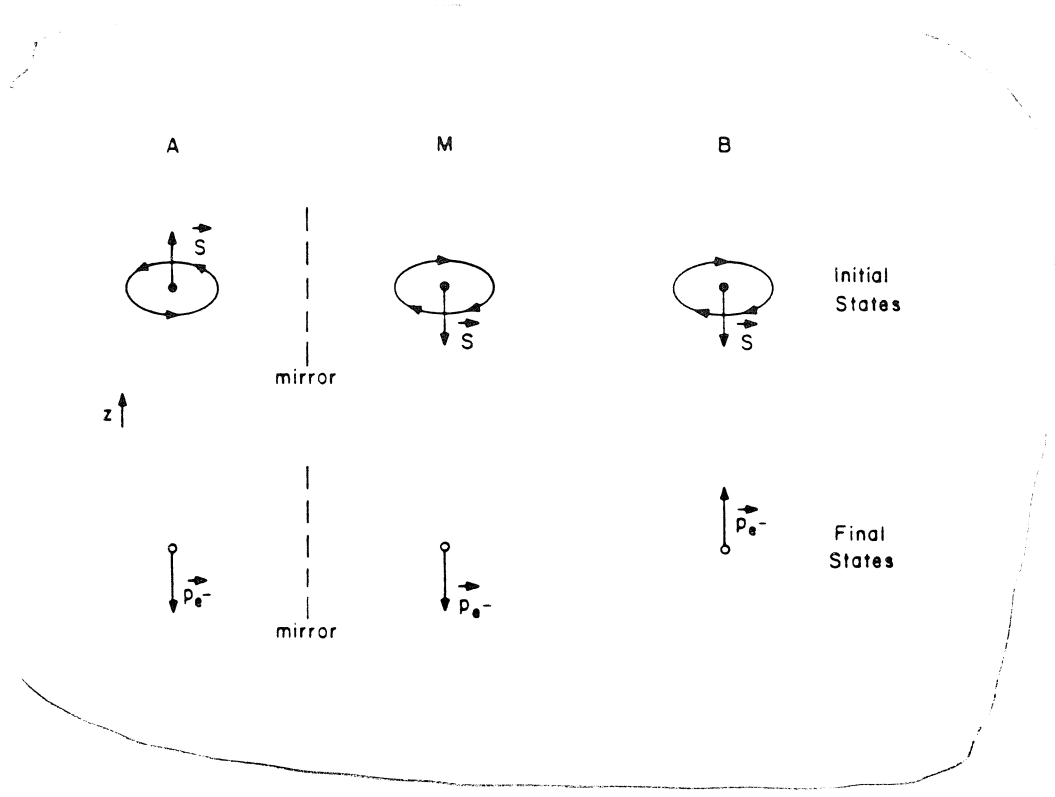
WEAK INTERACTION

NON PARITY

$$U(t) |\psi(0)\rangle \xleftrightarrow{\text{NON PARITY}} U(t) \pi |\psi(0)\rangle$$



flies out in a direction opposite to the spin of ${}^{60}\text{Co}$

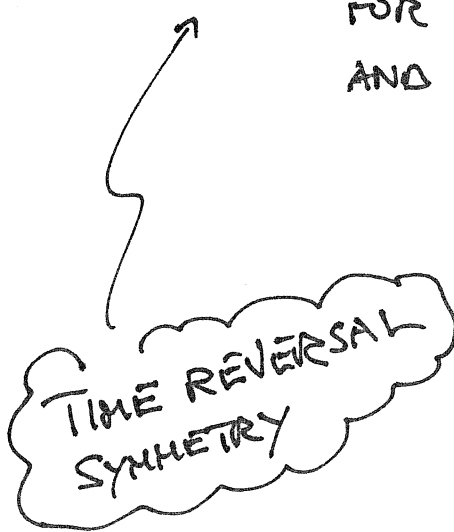


Lee & Yang (1956) theory
 Wu et al. (1957) experiment

STUDY CHAPTER 4
IN SAKURAI CAREFULLY !

4.1 - 4.3 EASY READ !

4.4 SEE THE COURSE HOME PAGE
FOR READING INSTRUCTIONS
AND SUPPLEMENTARY NOTES



TIME REVERSAL
SYMMETRY