Development of Many-Body Perturbation Theory

How to combine with Quantum ElectroDynamics?

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WHY?

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Seemingly incompatible

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 $\Psi(t_1 x_1, t_2 x_2..)$

Compromise: Equal-time approximation

Seemingly incompatible MBPT is based upon quantum mechanics with a single time

 $\Psi({m t},{m x}_1,{m x}_2..)$

QED is based on relativistic field theory with individual times

 $\Psi(\boldsymbol{t_1x_1}, \boldsymbol{t_2x_2}..)$

Compromise: Equal-time approximation

Very small effect in atomic/molecular physics

To treat correlation by QED is highly inefficient. Works for highly charged ions with small correlation relative to QED

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For chemical systems the situation is the reversed Correlated wave function natural starting point

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For chemical systems the situation is the reversed Correlated wave function natural starting point

First-order QED effects can be added to the ENERGY OFTEN INSUFFICIENT

QED perturbations are time- or energy dependent

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Requires time/energy-dependent perturbation theory

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Leads to a procedure where the QED effects are included perturbatively, mixed with the electron correlation - not added on at the final end

Outline

Review of standard methods

- Rayleigh-Schrödinger perturbation. Linked-diagram theorem. Bloch equation
- Relativistic MBPT. No-Virtual-Pair Approx.
- All-order methods. Coupled-cluster theory
- Methods for QED calculations

Outline

Beyond standard methods?

- Covariant Evolution Operator method
- Combination of MBPT and QED
- Numerical illustration: He-like systems
- Possible application to larger systems.
 Coupled-Cluster QED

MBPT Calculations Standard non-relativistic MBPT $H\Psi^{\alpha} = E^{\alpha}\Psi^{\alpha} \quad (\alpha = 1 \cdots d)$ $\Psi^{\alpha} = \mathbf{\Omega}\Psi^{\alpha}_{0} \quad (\alpha = 1 \cdots d) \quad \mathbf{\Omega}$ wave operator $\Psi^{\alpha}_{0} = P\Psi^{\alpha}$ Intermediate normalization **MBPT Calculations** Standard non-relativistic MBPT $H\Psi^{\alpha} = E^{\alpha}\Psi^{\alpha} \quad (\alpha = 1 \cdots d)$ $\Psi^{\alpha} = \mathbf{\Omega}\Psi^{\alpha}_{0} \quad (\alpha = 1 \cdots d) \quad \mathbf{\Omega}$ wave operator $\Psi^{\alpha}_{0} = P\Psi^{\alpha}$ Intermediate normalization

Linked diagram theorem

Graphical representation:Unlinked diagrams cancel (Brueckner 1955, Goldstone 1957, Brandow 1963, Mukherjee 1986) **MBPT** Calculations Standard non-relativistic MBPT $H\Psi^{\alpha} = E^{\alpha}\Psi^{\alpha} \quad (\alpha = 1 \cdots d)$ $\Psi^{\alpha} = \Omega \Psi_{0}^{\alpha}$ $(\alpha = 1 \cdots d)$ Ω wave operator $\Psi_0^{\alpha} = P\Psi^{\alpha}$ Intermediate normalization Bloch equation (Bloch 1958, IL 1974)

$$[\Omega, H_0]P = Q(V\Omega - \Omega W)_{\text{linked}}P$$

$$-\Omega W$$
 Model-space contribution
$$W = PV \Omega P$$
 Effective interaction

 $\left[\boldsymbol{\Omega}, H_0\right] P = Q \left(V \boldsymbol{\Omega} - \boldsymbol{\Omega} W \right)_{\text{linked}} P$ Second order : $\boldsymbol{\Omega}^{(2)} P_{\mathcal{E}} = \Gamma_Q \left(V \boldsymbol{\Omega}^{(1)} - \boldsymbol{\Omega}^{(1)} P_{\mathcal{E}'} W^{(1)} \right) P_{\mathcal{E}}$

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$$\left(\Omega^{(1)}(\mathcal{E}) - \Omega^{(1)}(\mathcal{E}')\right) \frac{P_{\mathcal{E}'}VP_{\mathcal{E}}}{\mathcal{E} - \mathcal{E}'} = \frac{\delta\Omega^{(1)}}{\delta\mathcal{E}} P_{\mathcal{E}'}VP_{\mathcal{E}}$$



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$$\Omega = \Gamma_Q \left[V\Omega - \Omega W\right] \text{holded}$$

Relativistic MBPT

(Breit 1931, Brown-Ravenhall 1951, Sucher 1980) Dirac-Coulomb-Breit Approximation

$$H = \mathbf{\Lambda}_{+} \Big[\sum_{i=1}^{N} h_{\mathbf{D}}(i) + \sum_{i < j}^{N} \frac{e^{2}}{4\pi r_{ij}} + H_{B} \Big] \mathbf{\Lambda}_{+}$$

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$$H_{B} = -\frac{e^{2}}{8\pi} \sum_{i < j} \Big[\frac{\alpha_{i} \cdot \alpha_{j}}{r_{ij}} + \frac{(\alpha_{i} \cdot r_{ij})(\alpha_{j} \cdot r_{ij})}{r_{ij}^{3}} \Big]$$

Instantaneous Breit interaction

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Instantaneous Breit interaction

No-(virtual)-pair approximation (NVPA)

QED effects

Effects beyond NVPA - Energy dependent Order α^3 and higher

- Retardation
- Virtual pairs Non-radiative
- Radiative effects (Lamb shift etc.)





Radiative effects

All-order methods

 $\Omega = 1 + \Omega_1 + \Omega_2 + \cdots$

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All-order pair function (Notre Dame, Gothenburg ...)



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Coulomb interaction

Folded

Internal vertical lines: electron propagators w pos. and neg. electron states Contains electron pair correlation to arbitrary order
Coupled-cluster approach

Exponential Ansatz (Coster 1958, Kümmel 1972, Čižek 1965) Closed shells (single reference)

$$\Omega = e^{T} = 1 + T + \frac{1}{2}T^{2} + \frac{1}{3!}T^{3} + \cdots$$

All diagrams connected

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Bloch equation

$$(E_0 - H_0)T = (V\Omega P - \Omega W)_{\text{com}}$$

Coupled-cluster approach Normal-ordered exponential Ansatz (IL1978, Mukherjee 1995, 97)

Open shells (multiple reference)

$$\Omega = \{e^T\} = 1 + T + \frac{1}{2}\{T^2\} + \frac{1}{3!}\{T^3\} + \cdots$$

Eliminates spurious contractions between open-shell lines

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Bloch equation

$$[T, H_0]P = Q(V\Omega - \Omega W)_{\text{conn}}$$

Coupled-cluster approach Open shells



Numerical, closed shells

BH₃ molecule (Shavitt et al. 1972)

Ecxitations	Total	Connected	Disconnected
One-body	0.1	0.1	
Two-body	97.2	97.2	≤ 0.1
Three-body	0.8	0.8	≤ 0.01
Four-body	1.9	≤ 0.01	1.9

Standard methods

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Standard methods treat energy-independent (Coulomb) interactions and energy-dependent QED perturbations separately

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Employing time-dependent perturbation theory enables us to combine MBPT and QED

Makes it possible to treat QED and correlation perturbatively on the same footing





A procedure for time-dependent perturbation theory has been developed based upon the **Covariant Evolution Operator (CEO)** method **Time-dependent perturbation theory** Covariant Evolution Operator (CEO) **Time-dependent perturbation theory Covariant Evolution Operator (CEO)** The single-particle **Green-s function** can be defined (in Heisenberg representation, *T* Wick time ordering)

 $G(t, t_0) = \frac{\left\langle 0_{\rm H} \middle| T[\hat{\psi}_{\rm H}(x)\hat{\psi}_{\rm H}^{\dagger}(x_0)] \middle| 0_{\rm H} \right\rangle}{\left\langle 0_{\rm H} \middle| 0_{\rm H} \right\rangle}$

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The single-particle **CEO** analogously

 $U_{\text{Cov}}(t,t_0) = \iint \mathrm{d}^3 \boldsymbol{x} \, \mathrm{d}^3 \boldsymbol{x}_0 \, \hat{\boldsymbol{\psi}}^{\dagger}(\boldsymbol{x}) \langle 0_{\mathrm{H}} \big| T[\hat{\psi}_{\mathrm{H}}(\boldsymbol{x})\hat{\psi}_{\mathrm{H}}^{\dagger}(\boldsymbol{x}_0)] \big| 0_{\mathrm{H}} \rangle \hat{\boldsymbol{\psi}}(\boldsymbol{x}_0)$

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GF is a function CEO is an operator

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Green's operator

Covariant evolution operator represents the time evolution of relativistic wave function

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Evolution operator singular due to intermediate model-space states

Green's operator

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 $\Psi_{\text{Rel}}(t) = U_{\text{Cov}}(t, t_0) \Psi_{\text{Rel}}(t_0)$

Evolution operator singular due to intermediate model-space states

Regular part known as Green's operator

 $U_{\text{Cov}}(t,t_0)P = \mathcal{G}(t,t_0)PU_{\text{Cov}}(0,t_0)P$

The Green's operator acts as the time-dependent relativistic wave operator

$$\Psi^lpha(t) = U_{
m Cov}(t,t_0) \Psi^lpha(t_0)$$

$$\Psi^{lpha}(t) = U_{
m Cov}(t,-\infty)\Psi^{lpha}(-\infty)$$

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Compare std MBPT:
$$\Psi = \Omega \Psi_0$$

Time-dependent perturbation theory Expansion of the Green's operator $U(t)P_{\mathcal{E}} = \mathcal{G}(t) P_{\mathcal{E}'}U(0)P_{\mathcal{E}}$ $\mathcal{G}^{(2)}(t)P_{\mathcal{E}} = Q\Big(\mathcal{G}^{(1)}(t) U^{(1)}(0) - \mathcal{G}^{(1)}(t)P_{\mathcal{E}'}U^{(1)}(0)\Big)P_{\mathcal{E}}$ **Time-dependent perturbation theory** Expansion of the Green's operator $U(t)P_{\mathcal{E}} = \mathcal{G}(t) P_{\mathcal{E}'}U(0)P_{\mathcal{E}}$ $\mathcal{G}^{(2)}(t)P_{\mathcal{E}} = Q\left(\mathcal{G}^{(1)}(t) U^{(1)}(0) - \mathcal{G}^{(1)}(t)P_{\mathcal{E}'}U^{(1)}(0)\right)P_{\mathcal{E}}$

Compare std wave operator

$$\Omega P = \Gamma_Q \left(V \Omega - \Omega P W \right) P$$
$$\mathbf{\Omega}^{(2)} P_{\mathcal{E}} = \Gamma_Q \left(V \mathbf{\Omega}^{(1)} - \mathbf{\Omega}^{(1)} P_{\mathcal{E}'} W^{(1)} \right) P_{\mathcal{E}}$$
Folded diagrams



 $\mathbf{\Omega}^{(2)} = \Gamma_Q V \mathbf{\Omega}^{(1)} + \frac{\delta \mathbf{\Omega}^{(1)}}{\delta \mathcal{E}} W^{(1)} = \Gamma_Q V \mathbf{\Omega}^{(1)} - \Gamma_Q \mathbf{\Omega}^{(1)} W^{(1)}$

Energy-independent perturbations

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Energy-independent perturbations

$$\mathcal{G}^{(2)} = \Gamma_Q V \mathcal{G}^{(1)} - \Gamma_Q \mathcal{G}^{(1)} W^{(1)} + \Gamma_Q \frac{\delta V}{\delta \mathcal{E}} W^{(1)}$$

Energy-dependent perturbations

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Energy-dependent perturbations

Bloch eqn for time-dependent perturbation theory

$$ig[\mathcal{G}, H_0 ig] = V \mathcal{G} - \mathcal{G} W + ig[rac{\delta^* \mathcal{G}}{\delta \mathcal{E}}, H_0 ig] W$$

* derivation restricted to last interaction

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ight]=V\Omega-\Omega W$$

 $(W = PV\Omega P)$

Procedures completely compatible

How to evaluate the QED effects perturbatively?

Gell-Mann-Low theorem

$$ig|oldsymbol{H}_Dig|\Psiig
angle=ig(oldsymbol{H}_0+V_C+v_Tig)ig|\Psiig
angle=oldsymbol{E}ig|\Psiig
angle$$

Gell-Mann-Low theorem

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$$V = V_C + v_T$$

Coulomb + Transverse

Gell-Mann-Low theorem

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Coulomb + Transverse

 $V_C=rac{e^2}{4\pi r_{12}}$

$$v_T = -\int \mathrm{d}^3x\, \hat{\psi}(x)^\dagger e c lpha^\mu A_\mu(x) \hat{\psi}(x)$$



Photonic Fock space



Photonic Fock space

Perturbation should be time independent! for Gell-Mann-Low to be valid

Single-photon exchange requires two perturbations



Single-photon exchange requires two perturbations



Single-photon exchange requires two perturbations

$$r = \int_{t}^{r} \int_{0}^{\infty} d\kappa f(\kappa) \frac{1}{\mathcal{E} - \varepsilon_r - \varepsilon_u - \kappa} |tu\rangle$$
$$f(\kappa) = \sum_{t} V^{l}(\kappa r_1) \cdot V^{l}(\kappa r_2)$$

TWO energy-independent perturbations. The energy dependence is given by the energy denominator. GML valid

$$|H_D|\Psi_l\rangle = \left(H_0 + V_C + V^l\right)|\Psi_l\rangle = E|\Psi_l\rangle$$

















Time-dependent perturbation theory Calculating derivative of retarded interaction



$$V = \frac{f(k)}{\mathcal{E} - \varepsilon_r - \varepsilon_q - k}$$

Derivative by second energy denominator



Self-energy regularization

Dimensional regularization in Coulomb gauge most appropriate to use in MBPT/QED Never used before.



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Zero- and one-pot. terms evaluated using Adkins formulas Modified by J. Holmberg, PRA**84**, 062504 (2011)

Many-potential term obtained by evaluating the other terms with partial-wave expansion

Self-energy regularization

First dimensional regularization in Coulomb gauge Self-energy of hydrogen like ions

Z	Coulomb gauge	Feynman gauge
18	1.216901(3)	1.21690(1)
54	50.99727(2)	50.99731(8)
66	102.47119(3)	102.4713(1)
92	355.0430(1)	355.0432(2)

D. Hedendahl and J. Holmberg, Phys. Rev. A85, 012514 (2012)

Self-energy

Derivative of self-energy is singular Singularity cancelled by vertex correction due to Ward identity

$$\frac{\delta \Sigma}{\delta \mathcal{E}} = \Lambda_0$$

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$$\frac{\delta \Sigma}{\delta \mathcal{E}} = \Lambda_0$$

Both are charge divergent and have to be renormalized





















Numerical illustration

QED effects in He-like ions, grd state (eV)

Ζ	Two-photon	Combined QED-correlation BEYOND two-photon					
	Retarded	Inst. Breit	Retard. part	Virt.Pairs	Self-energy	Vertex corr.	
10	0.0033	0.0072	-0.0011	0.0002			
14	0.0080	0.0101	-0.0019	0.0004	0.0020		
18	0.0150	0.0154	-0.0027	0.0006	0.0030		
24	0.0305	0.0192	-0.0042	0.0009	0.0050		
30	0.052	0.0244	-0.0057	0.0013	0.0090		
42	0.112	0.0286	-0.0087	0.0019			
50		0.0320	-0.011	0.0024	0.0170		
66		0.0400	-0.015	0.0030			
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CCSD(T)

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Coupled-cluster CCSD



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Coupled-cluster CCSD



MBPT-QED

QED potentials



Coupled-cluster CCSD(T)



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The procedure has been tested for He-like ions

Application to more general systems might be possible by using the coupled-cluster approach

Spitinger Series on Atomic, Optical and Plasma Physics

Ingvar Lindgren

Relativistic Many-Body Theory

A New Field-Theoretical Approach

D Springer

Coworkers

Sten Salomonson Daniel Hedendahl Johan Holmberg Thank you!