# Development of Many-Body Perturbation Theory 

How to combine with Quantum ElectroDynamics?

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WHY?

## Fundamental problem

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Compromise: Equal-time approximation
Very small effect in atomic/molecular physics

## How can QED effects be included in a many-body problem in a systematic fashion?

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For chemical systems the situation is the reversed Correlated wave function natural starting point

First-order QED effects can be added to the ENERGY OFTEN INSUFFICIENT

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## To go further, QED effects should be included in the

 WAVE FUNCTIONQED perturbations are time- or energy dependent
Requires time/energy-dependent perturbation theory
Leads to a procedure where the QED effects are included perturbatively, mixed with the electron correlation - not added on at the final end

## Outline

## Review of standard methods

- Rayleigh-Schrödinger perturbation. Linked-diagram theorem. Bloch equation
- Relativistic MBPT. No-Virtual-Pair Approx.
- All-order methods. Coupled-cluster theory
- Methods for QED calculations


## Outline

## Beyond standard methods?

- Covariant Evolution Operator method
- Combination of MBPT and QED
- Numerical illustration: He-like systems
- Possible application to larger systems. Coupled-Cluster QED


## MBPT Calculations

## Standard non-relativistic MBPT

$$
\begin{gathered}
H \Psi^{\alpha}=E^{\alpha} \Psi^{\alpha} \quad(\alpha=1 \cdots d) \\
\Psi^{\alpha}=\Omega \Psi_{0}^{\alpha} \quad(\alpha=1 \cdots d) \quad \Omega \text { wave operator } \\
\Psi_{0}^{\alpha}=P \Psi^{\alpha} \quad \text { Intermediate normalization }
\end{gathered}
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## Linked diagram theorem

Graphical representation:Unlinked diagrams cancel
(Brueckner 1955, Goldstone 1957, Brandow 1963, Mukherjee 1986)

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Bloch equation (Bloch 1958, IL 1974)
$\left[\Omega, H_{0}\right] P=Q(V \Omega-\Omega W)_{\text {linked }} P$

$-\Omega W$
Model-space contribution
$W=P V \Omega P \quad$ Effective interaction

## Model-space contribution

$$
\left[\Omega, H_{0}\right] P=Q(V \Omega-\Omega W)_{\text {linked }} P
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Second order : $\Omega^{(2)} P_{\mathcal{E}}=\Gamma_{Q}\left(V \Omega^{(1)}-\Omega^{(1)} P_{\mathcal{E}} W^{(1)}\right) P_{\mathcal{E}}$

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$\left\{\begin{array}{l}1 \\ \Omega^{(1)}\left(\mathcal{E}^{\prime}\right) \\ 1\end{array}\right.$
1


Unlinked


Unlinked Linked

## Model-space contribution

Remainder: Folded diagrams


Folded

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$$
\left(\Omega^{(1)}(\mathcal{E})-\Omega^{(1)}\left(\mathcal{E}^{\prime}\right)\right) \frac{P_{\mathcal{E}^{\prime}} V P_{\mathcal{E}}}{\mathcal{E}-\mathcal{E}^{\prime}}=\frac{\delta \Omega^{(1)}}{\delta \mathcal{E}} P_{\mathcal{E}^{\prime}} V P_{\mathcal{E}}
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\Omega=\Gamma_{Q}[V \Omega-\Omega W]_{\text {linked }}
\end{gathered}
$$

## Relativistic MBPT

(Breit 1931, Brown-Ravenhall 1951, Sucher 1980)
Dirac-Coulomb-Breit Approximation

$$
H=\Lambda_{+}\left[\sum_{i=1}^{N} h_{D}(i)+\sum_{i<j}^{N} \frac{e^{2}}{4 \pi r_{i j}}+H_{B}\right] \Lambda_{+}
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Instantaneous Breit interaction

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## QED effects

Effects beyond NVPA - Energy dependent

## Order $\boldsymbol{\alpha}^{3}$ and higher

- Retardation
- Virtual pairs Non-radiative
- Radiative effects (Lamb shift etc.)



## All-order methods

$$
\Omega=1+\Omega_{1}+\Omega_{2}+\cdots
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----- Coulomb interaction
$\rightleftharpoons$ Folded
Internal vertical lines: electron propagators w pos. and neg. electron states
Contains electron pair correlation to arbitrary order

## Coupled-cluster approach

Exponential Ansatz (Coster 1958, Kümmel 1972, Čǐ̌ek 1965)
Closed shells (single reference)

$$
\Omega=\mathrm{e}^{T}=1+T+\frac{1}{2} T^{2}+\frac{1}{3!} T^{3}+\cdots
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## All diagrams connected

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All diagrams connected
Bloch equation

$$
\left(E_{0}-H_{0}\right) T=(V \Omega P-\Omega W)
$$

## Coupled-cluster approach

Normal-ordered exponential Ansatz
( IL1978, Mukherjee 1995, 97)
Open shells (multiple reference)

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Eliminates spurious contractions between open-shell lines

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## Bloch equation

$$
\left[T, H_{0}\right] P=Q(V \Omega-\Omega W)_{\mathrm{conn}}
$$

## Coupled-cluster approach

Open shells

$$
\begin{aligned}
& f=f(1)+\frac{f}{f}+\cdots+\frac{1}{F}
\end{aligned}
$$

## Numerical, closed shells

$\mathrm{BH}_{3}$ molecule (Shavitt et al. 1972)

| Ecxitations | Total | Connected | Disconnected |
| :---: | :---: | :---: | :---: |
| One-body | 0.1 | 0.1 |  |
| Two-body | 97.2 | 97.2 | $\leq 0.1$ |
| Three-body | 0.8 | 0.8 | $\leq 0.01$ |
| Four-body | 1.9 | $\leq 0.01$ | 1.9 |

## QED Calculations

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## Beyond standard methods

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Standard methods treat energy-independent (Coulomb) interactions and energy-dependent QED perturbations separately
This leaves out the combination effect
Employing time-dependent perturbation theory enables us to combine MBPT and QED

Makes it possible to treat QED and correlation perturbatively on the same footing

## Time-dependent perturbation theory



## Time-dependent perturbation theory



A procedure for time-dependent perturbation theory has been developed based upon the
Covariant Evolution Operator (CEO) method

## Time-dependent perturbation theory Covariant Evolution Operator (CEO)

## Time-dependent perturbation theory

 Covariant Evolution Operator (CEO)The single-particle Green-s function can be defined (in Heisenberg representation, $T$ Wick time ordering)

$$
G\left(t, t_{0}\right)=\frac{\left\langle 0_{\mathrm{H}}\right| T\left[\hat{\psi}_{\mathrm{H}}(x) \hat{\psi}_{\mathrm{H}}^{\dagger}\left(x_{0}\right)\right]\left|0_{\mathrm{H}}\right\rangle}{\left\langle 0_{\mathrm{H}} \mid 0_{\mathrm{H}}\right\rangle}
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$$

The single-particle CEO analogously

$$
U_{\mathrm{Cov}}\left(t, t_{0}\right)=\iint \mathrm{d}^{3} \boldsymbol{x} \mathrm{~d}^{3} \boldsymbol{x}_{0} \hat{\psi}^{\dagger}(x)\left\langle 0_{\mathrm{H}}\right| T\left[\hat{\psi}_{\mathrm{H}}(x) \hat{\psi}_{\mathrm{H}}^{\dagger}\left(x_{0}\right)\right]\left|0_{\mathrm{H}}\right\rangle \hat{\psi}\left(x_{0}\right)
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GF is a function CEO is an operator

## Time-dependent perturbation theory

Green's operator
Covariant evolution operator represents the time evolution of relativistic wave function

$$
\Psi_{\operatorname{Rel}}(t)=U_{\mathrm{Cov}}\left(t, t_{0}\right) \Psi_{\mathrm{Rel}}\left(t_{0}\right)
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Evolution operator singular due to intermediate model-space states

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Evolution operator singular due to intermediate model-space states

Regular part known as Green's operator

$$
U_{\mathrm{Cov}}\left(t, t_{0}\right) P=\mathcal{G}\left(t, t_{0}\right) P U_{\mathrm{Cov}}\left(0, t_{0}\right) P
$$

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The Green's operator acts as the time-dependent relativistic wave operator

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$$

Compare std MBPT: $\Psi=\Omega \Psi_{0}$

## Time-dependent perturbation theory

Expansion of the Green's operator

$$
\begin{gathered}
U(t) P_{\mathcal{E}}=\mathcal{G}(t) P_{\mathcal{E}} U(0) P_{\mathcal{E}} \\
\mathcal{G}^{(2)}(t) P_{\mathcal{E}}=Q\left(\mathcal{G}^{(1)}(t) U^{(1)}(0)-\mathcal{G}^{(1)}(t) P_{\mathcal{E}} U^{(1)}(0)\right) P_{\mathcal{E}}
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\end{gathered}
$$

Compare std wave operator

$$
\begin{gathered}
\Omega P=\Gamma_{Q}(V \Omega-\Omega P W) P \\
\Omega^{(2)} P_{\mathcal{E}}=\Gamma_{Q}\left(V \Omega^{(1)}-\Omega^{(1)} P_{\mathcal{E}} W^{(1)}\right) P_{\mathcal{E}}
\end{gathered}
$$

## Time-dependent perturbation theory

Folded diagrams


Folded

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\Omega^{(2)}=\Gamma_{Q} V \Omega^{(1)}+\frac{\delta \Omega^{(1)}}{\delta \mathcal{E}} W^{(1)}=\Gamma_{Q} V \Omega^{(1)}-\Gamma_{Q} \Omega^{(1)} W^{(1)}
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Energy-independent perturbations

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Energy-dependent perturbations

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## Bloch eqn for time-dependent perturbation theory

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\left[\mathcal{G}, \boldsymbol{H}_{0}\right]=\boldsymbol{V} \mathcal{G}-\mathcal{G} W+\left[\frac{\delta^{*} \mathcal{G}}{\delta \mathcal{E}}, \boldsymbol{H}_{0}\right] \boldsymbol{W}
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* derivation restricted to last interaction


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## Procedures completely compatible

# Time-dependent perturbation theory 

## How to evaluate the QED effects perturbatively?

## Time-dependent perturbation theory

## Gell-Mann-Low theorem

$$
H_{D}|\Psi\rangle=\left(H_{0}+V_{C}+v_{T}\right)|\Psi\rangle=E|\Psi\rangle
$$

## Time-dependent perturbation theory

## Gell-Mann-Low theorem

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$$
V=V_{C}+v_{T}
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Coulomb + Transverse

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Coulomb + Transverse

$$
V_{C}=\frac{e^{2}}{4 \pi r_{12}}
$$

$$
v_{T}=-\int \mathrm{d}^{3} x \hat{\psi}(x)^{\dagger} e c \alpha^{\mu} A_{\mu}(x) \hat{\psi}(x)
$$

## Time-dependent perturbation theory



## Photonic Fock space

## Time-dependent perturbation theory



## Photonic Fock space

## Perturbation should be time independent! for Gell-Mann-Low to be valid

## Time-dependent perturbation theory

Single-photon exchange requires two perturbations

$$
\langle r s| V_{\mathrm{sp}}(\mathcal{E})|t u\rangle=\langle r s| \int_{0}^{\infty} \mathrm{d} \kappa f(\kappa) \frac{1}{\mathcal{E}-\varepsilon_{r}-\varepsilon_{u}-\kappa}|t u\rangle
$$

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Single-photon exchange requires two perturbations

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\begin{gathered}
\langle r s| V_{\mathrm{sp}}(\mathcal{E})|t u\rangle=\langle r s| \int_{0}^{\infty} \mathrm{d} \kappa f(\kappa) \frac{1}{\mathcal{E}-\varepsilon_{r}-\varepsilon_{u}-\kappa}|t u\rangle \\
f(\kappa)=\sum_{l} V^{l}\left(\kappa \boldsymbol{r}_{1}\right) \cdot V^{l}\left(\kappa \boldsymbol{r}_{2}\right)
\end{gathered}
$$

TWO energy-independent perturbations. The energy dependence is given by the energy denominator. GML valid

$$
\boldsymbol{H}_{D}\left|\Psi_{l}\right\rangle=\left(\boldsymbol{H}_{0}+\boldsymbol{V}_{C}+V^{l}\right)\left|\Psi_{l}\right\rangle=\boldsymbol{E}\left|\Psi_{l}\right\rangle
$$

## Time-dependent perturbation theory

Iteration of time-independent perturbations


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## Time-dependent perturbation theory

Calculating derivative of retarded interaction


$$
V=\frac{f(k)}{\varepsilon-\varepsilon_{r}-\varepsilon_{q}-k}
$$

Derivative by second energy denominator


## Self-energy regularization

Dimensional regularization in Coulomb gauge most appropriate to use in MBPT/QED

Never used before.


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Zero- and one-pot. terms evaluated using Adkins formulas
Modified by J. Holmberg, PRA84, 062504 (2011)
Many-potential term obtained by evaluating the other terms with partial-wave expansion

## Self-energy regularization

First dimensional regularization in Coulomb gauge Self-energy of hydrogen like ions

| $\mathbb{Z}$ | Coulomb gauge | Feynman gauge |
| :---: | :---: | :---: |
| 18 | $1.216901(3)$ | $1.21690(1)$ |
| 54 | $50.99727(2)$ | $50.99731(8)$ |
| 66 | $102.47119(3)$ | $102.4713(1)$ |
| 92 | $355.0430(1)$ | $355.0432(2)$ |

D. Hedendahl and J. Holmberg, Phys. Rev. A85, 012514 (2012)

## Self-energy

## Derivative of self-energy is singular

 Singularity cancelled by vertex correction due to Ward identity$$
\frac{\delta \Sigma}{\delta \mathcal{E}}=\Lambda_{0}
$$

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Both are charge divergent and have to be renormalized

## Time-dependent perturbation theory

Iteration of time-dependent perturbations

## QED Pair function



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## Numerical illustration

## QED effects in He-like ions, grd state (eV)

| Z | Two-photon | Combined QED-correlation BEYOND two-photon |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | Retarded | Inst. Breit | Retard. part | Virt.Pairs | Self-energy | Vertex corr. |
| 10 | 0.0033 | 0.0072 | -0.0011 | 0.0002 |  |  |
| 14 | 0.0080 | 0.0101 | -0.0019 | 0.0004 | 0.0020 |  |
| 18 | 0.0150 | 0.0154 | -0.0027 | 0.0006 | 0.0030 |  |
| 24 | 0.0305 | 0.0192 | -0.0042 | 0.0009 | 0.0050 |  |
| 30 | 0.052 | 0.0244 | -0.0057 | 0.0013 | 0.0090 |  |
| 42 | 0.112 | 0.0286 | -0.0087 | 0.0019 |  |  |
| 50 |  | 0.0320 | -0.011 | 0.0024 | 0.0170 |  |
| 66 |  | 0.0400 | -0.015 | 0.0030 |  |  |

## Can this be applied to larger systems???

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Might have the potential, since the standard and QED procedures are completely compatible

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$$
\operatorname{CCSD}(\mathrm{T})
$$





## Coupled-cluster CCSD

$$
\begin{aligned}
& t=f-x+\frac{f}{f}+\underset{f}{f}+\frac{1}{f} \mathbb{L}+\cdots+\frac{t}{F} \\
& H=H+H+H+H_{+}^{+}+{ }_{+}^{+} \\
& \text {- }
\end{aligned}
$$

## Coupled-cluster CCSD

$$
\begin{aligned}
& \mathrm{H} \cdot \mathrm{H}+\mathrm{H} \cdot \mathrm{H}+\mathrm{H} \cdot \mathrm{H} \\
& \text { fol|fori| } \cdots+1 \cdot H \cdot+\#
\end{aligned}
$$

## MBPT-QED

## QED potentials

$$
\begin{aligned}
& Y=+x+1+H_{0} \\
& H=H+1+H+N+H+W+\cdots
\end{aligned}
$$

## Coupled-cluster CCSD(T)

$$
\begin{aligned}
& t=1+1+10+\frac{1}{1}+\cdots+\frac{1}{7} \\
& 1+1+\frac{1}{1}+\frac{1}{1}+\frac{1}{1}+\frac{1}{1}+ \\
& +0|+0|+\cdots+\frac{1}{1}+\frac{1}{1}+\frac{1}{1}
\end{aligned}
$$

## Summary

A procedure has been developed for combining MBPT and QED, based upon the covariant-evolution operator method

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The compatibility between this procedure and standard MBPT implies that the QED effects can be inserted to the degree that is relevant
The procedure has been tested for He-like ions
Application to more general systems might be possible by using the coupled-cluster approach


Ingurlindeten
Relativistic Many-Body Theory

AMew Fidd-hevelralAprosth

## Coworkers

## Sten Salomonson Daniel Hedendahl Johan Holmberg

## Thank you!

