

Lectures

Exercises/Home work problems

Ordinary differential equations

Linear dynamics

E1

Non-linear dynamics

E2

Molecular dynamics

H1a/H1b

Stochastic methods

Monte Carlo integration

E3

Metropolis algorithm

H2a/H2b

Brownian dynamics

E4

Partial differential equations

Quantum structure

E5

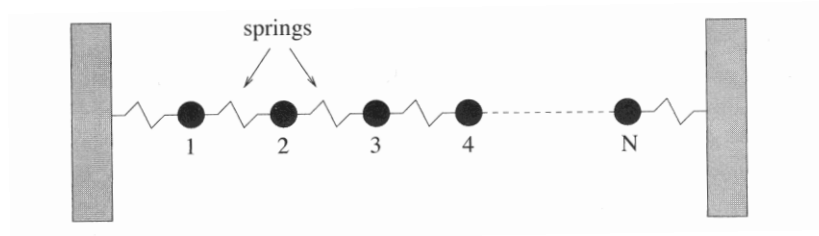
H3a

Quantum dynamics

H3b

Fermi-Pasta-Ulam model

One-dimensional chain of masses connected with non-linear springs.



α -model

$$H = \sum_{i=1}^N \frac{p_i^2}{2m} + \sum_{i=0}^N \left[\frac{\kappa}{2} (u_{i+1} - u_i)^2 + \frac{\alpha}{3} (u_{i+1} - u_i)^3 \right]$$

Fermi-Pasta-Ulam problem

After the second world war Fermi *et al.* became interested in using the newly developed computing machines to test physical ideas, to perform "computer experiments" or "computer simulations".

They wanted to check the prediction of statistical mechanics on the thermalization of solids, which relies on the ergodic hypothesis.

Print 013-018
Convert to decimals
& print 210-228

Automatic operation
Start 170

Summary routine (start 27E)
a) Input sum
b) 3A0ICFFSFE
and 26: FFFFE
and 26: FFFFE
constants

Conversion of input to hexadecimal (start 317)
Punch input (start 330)

Printout of Fermi's program. It is written in machine language.
from H. L. Anderson, J. Stat. Phys. 43, 731 (1986)

Ergodic hypothesis

Time average = Ensemble average

Consider a closed system

Classical mechanics

The system follows a deterministic time-evolution where the energy is conserved. The time-evolution is reversible with no direction in time.

Statistical mechanics

The system will move to a state of thermal equilibrium, where the entropy is maximized. The time-evolution is irreversible with a clear direction of time.

Equipartition theorem

In classical statistical mechanics the mean value of each independent quadratic term in the energy is equal to $\frac{k_B T}{2}$.

Mean value

$$\langle A \rangle = \frac{\int dx_1 \dots dp_N A(x_1, \dots, p_N) \exp[-E(x_1, \dots, dp_N)/k_B T]}{\int dx_1 \dots dp_N \exp[-E(x_1, \dots, dp_N)/k_B T]}$$

Consider

$$E(x_1, \dots, p_N) = ax_i^2 + E'(x_1, \dots, x_{i-1}, x_{i+1}, \dots, p_N)$$

which implies that

$$\langle ax_i^2 \rangle = \frac{k_B T}{2}$$

Equipartition theorem

In classical statistical mechanics the mean value of each independent quadratic term in the energy is equal to $\frac{k_B T}{2}$.

Energy of an ideal gas particle

$$E = \frac{\mathbf{p}^2}{2m} = \frac{p_x^2}{2m} + \frac{p_y^2}{2m} + \frac{p_z^2}{2m}$$
$$\langle E \rangle = \frac{3k_B T}{2}$$

System of interacting particles

$$E = \sum_{i=1}^N \frac{\mathbf{p}_i^2}{2m} + V(\mathbf{r}_1, \dots, \mathbf{r}_N)$$
$$\langle E_{kin} \rangle = N \frac{3k_B T}{2}$$

Harmonic oscillator in 1 dim

$$E = \frac{p^2}{2m} + \frac{k}{2} x^2$$
$$\langle E \rangle = \frac{k_B T}{2} + \frac{k_B T}{2} = k_B T$$

System of independent harmonic oscillators in 1 dim

$$E = \sum_{i=1}^N E_i = \sum_{i=1}^N \left[\frac{p_i^2}{2m} + m \frac{k}{2} x_i^2 \right]$$
$$\langle E_i \rangle = \left\langle \frac{p_i^2}{2m} + \frac{k}{2} x_i^2 \right\rangle = k_B T \quad \forall i$$
$$\langle E \rangle = N k_B T$$

Harmonic system

Hamiltonian

$$H = \sum_{i=1}^N \frac{p_i^2}{2m} + \sum_{i=0}^N \frac{\kappa}{2} (u_{i+1} - u_i)^2$$

Modal representation - normal modes

$$Q_k = \sqrt{\frac{2}{N+1}} \sum_{i=1}^N \sqrt{m} u_i \sin\left(\frac{ik\pi}{N+1}\right)$$

$$u_i = \sqrt{\frac{2}{N+1}} \sum_{k=1}^N \frac{Q_k}{\sqrt{m}} \sin\left(\frac{ik\pi}{N+1}\right)$$

and

$$P_k = \sqrt{\frac{2}{N+1}} \sum_{i=1}^N \sqrt{m} v_i \sin\left(\frac{ik\pi}{N+1}\right)$$

$$v_i = \sqrt{\frac{2}{N+1}} \sum_{k=1}^N \frac{P_k}{\sqrt{m}} \sin\left(\frac{ik\pi}{N+1}\right)$$

for $k = 1, \dots, N$ and $i = 1, \dots, N$.

Hamiltonian

$$H = \frac{1}{2} \sum_{k=1}^N [P_k^2 + \omega_k^2 Q_k^2]$$

with

$$\omega_k = 2\sqrt{\frac{\kappa}{m}} \sin\frac{k\pi}{2(N+1)}$$

i.e. independent harmonic oscillators.

Assuming equipartition

$$\left\langle \frac{1}{2} P_k^2 \right\rangle = \frac{k_B T}{2} \quad \forall k$$

$$\left\langle \frac{1}{2} \omega_k^2 Q_k^2 \right\rangle = \frac{k_B T}{2} \quad \forall k$$

and

$$\langle E_k \rangle = \left\langle \frac{1}{2} (P_k^2 + \omega_k^2 Q_k^2) \right\rangle = k_B T \quad \forall k$$

Fermi-Pasta-Ulam model

$$H = \sum_{i=1}^N \frac{p_i^2}{2m} + \sum_{i=0}^N \left[\frac{\kappa}{2} (u_{i+1} - u_i)^2 + \frac{\alpha}{3} (u_{i+1} - u_i)^3 \right]$$

$\alpha = 0$

Linear system.

can be described by normal modes

no energy exchange between the modes

no equipartition of the energy among the normal modes

$\alpha \neq 0$

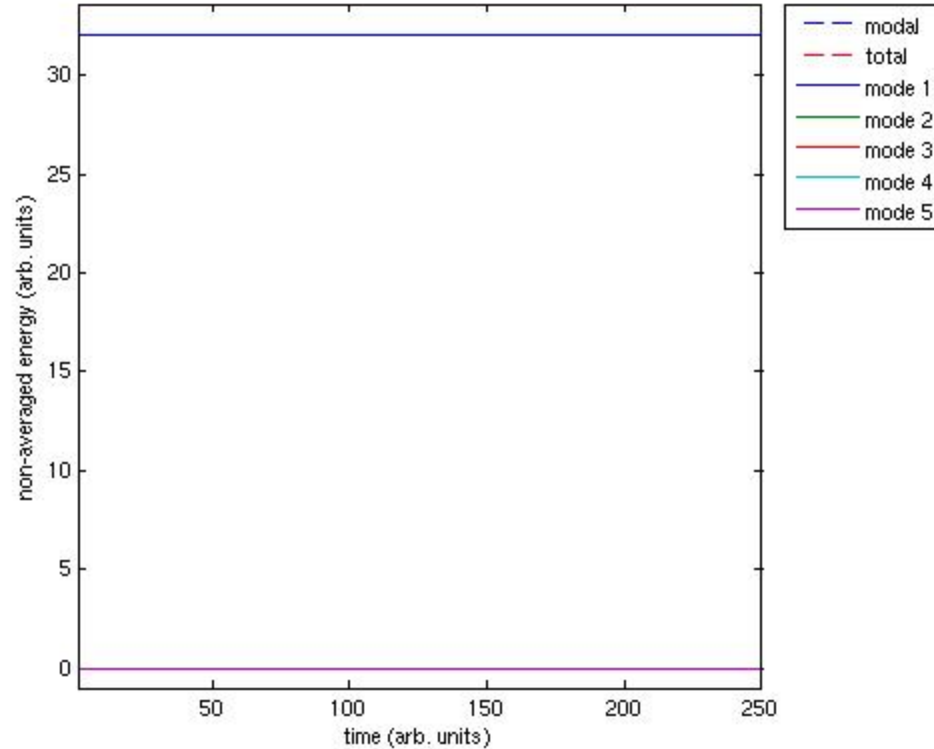
Non-linear system

can approximately be described by normal modes (weak non-linearity)

energy exchange between the modes

equipartition of the energy among the normal modes **?**

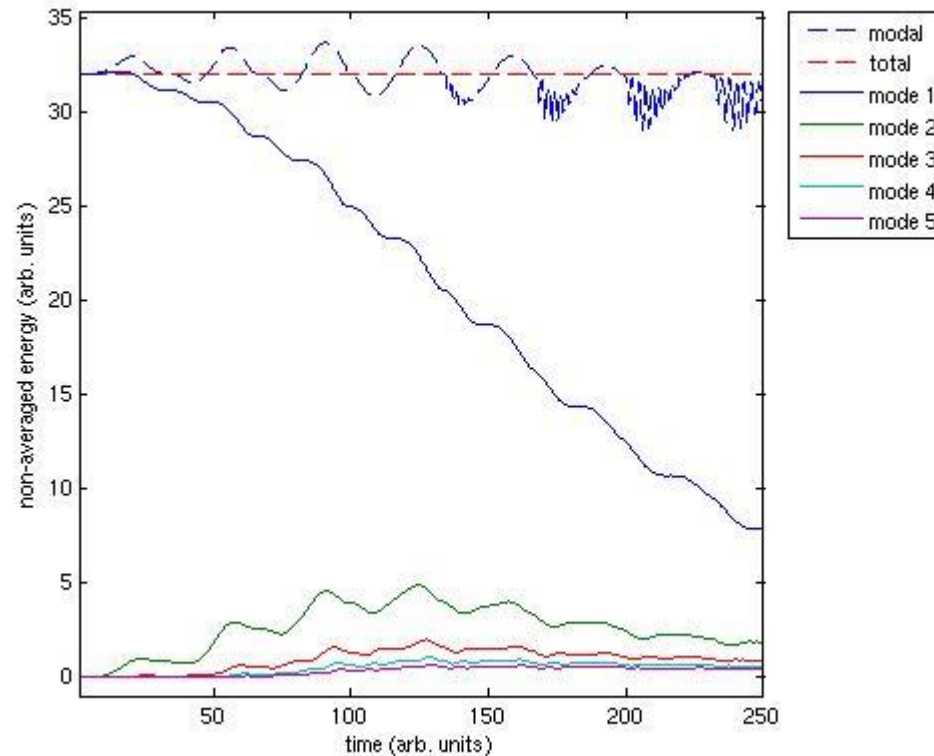
Harmonic system – numerical result



No equipartition of the energy

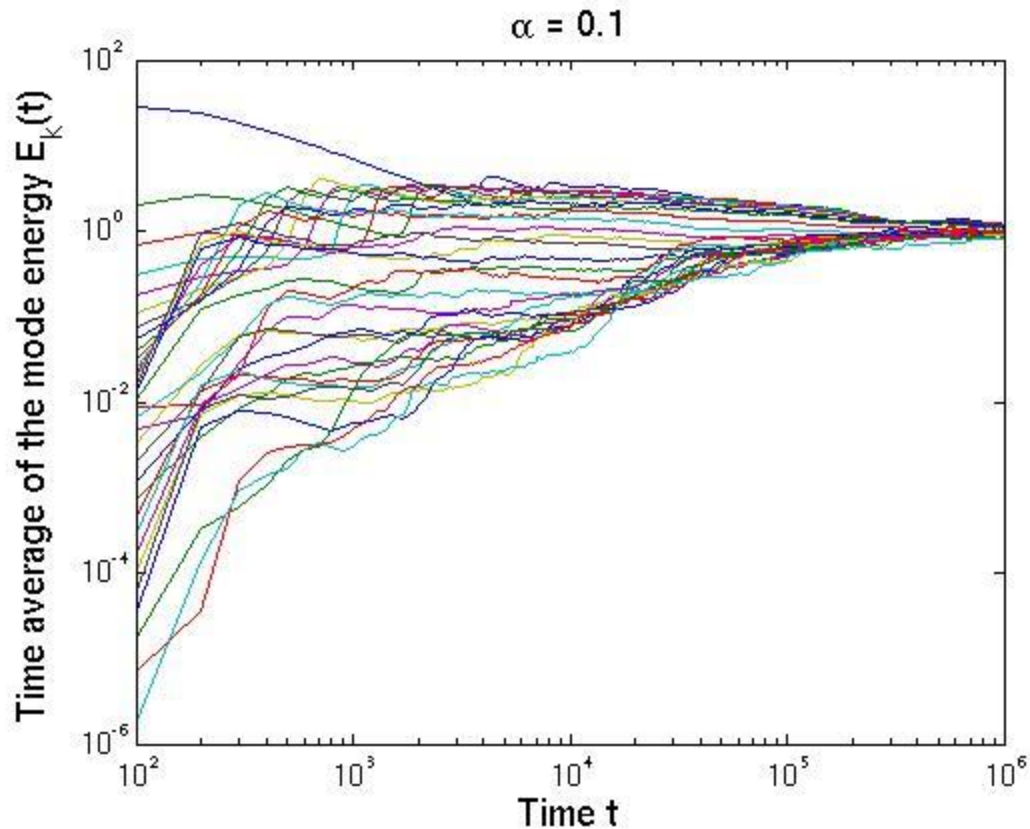
Anharmonic system – numerical result

$$\alpha = 0.1$$



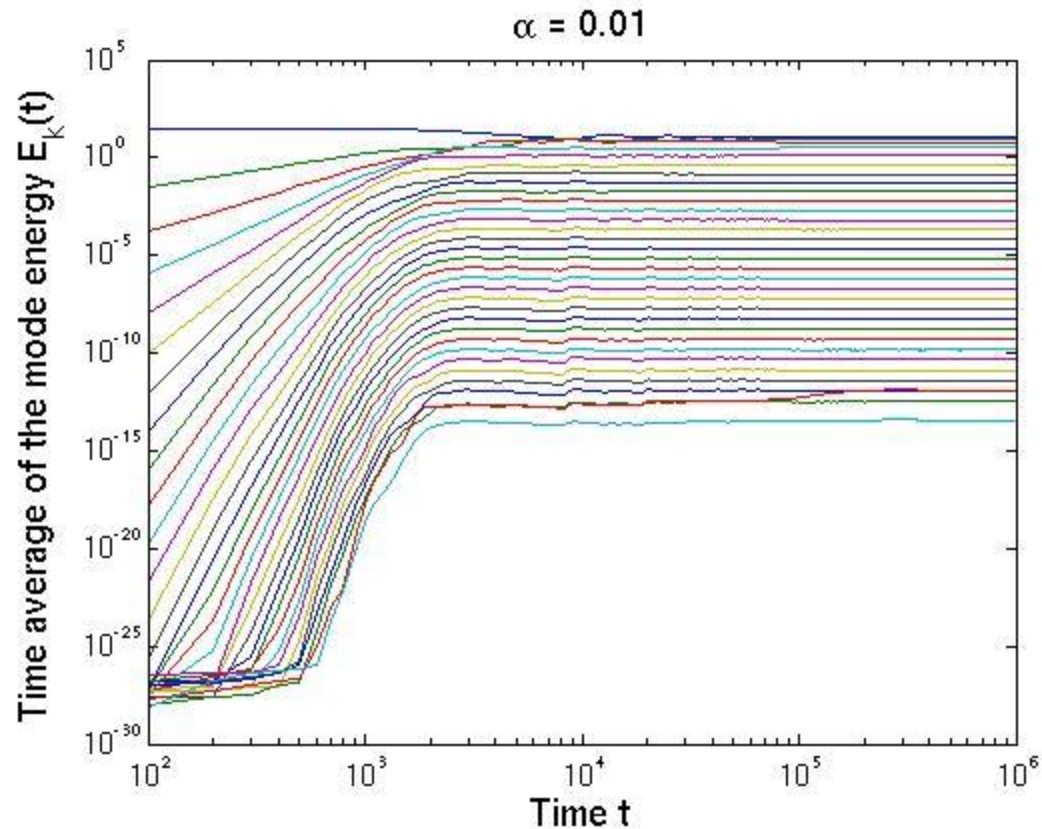
Equipartition of the energy ?

Anharmonic system – numerical result



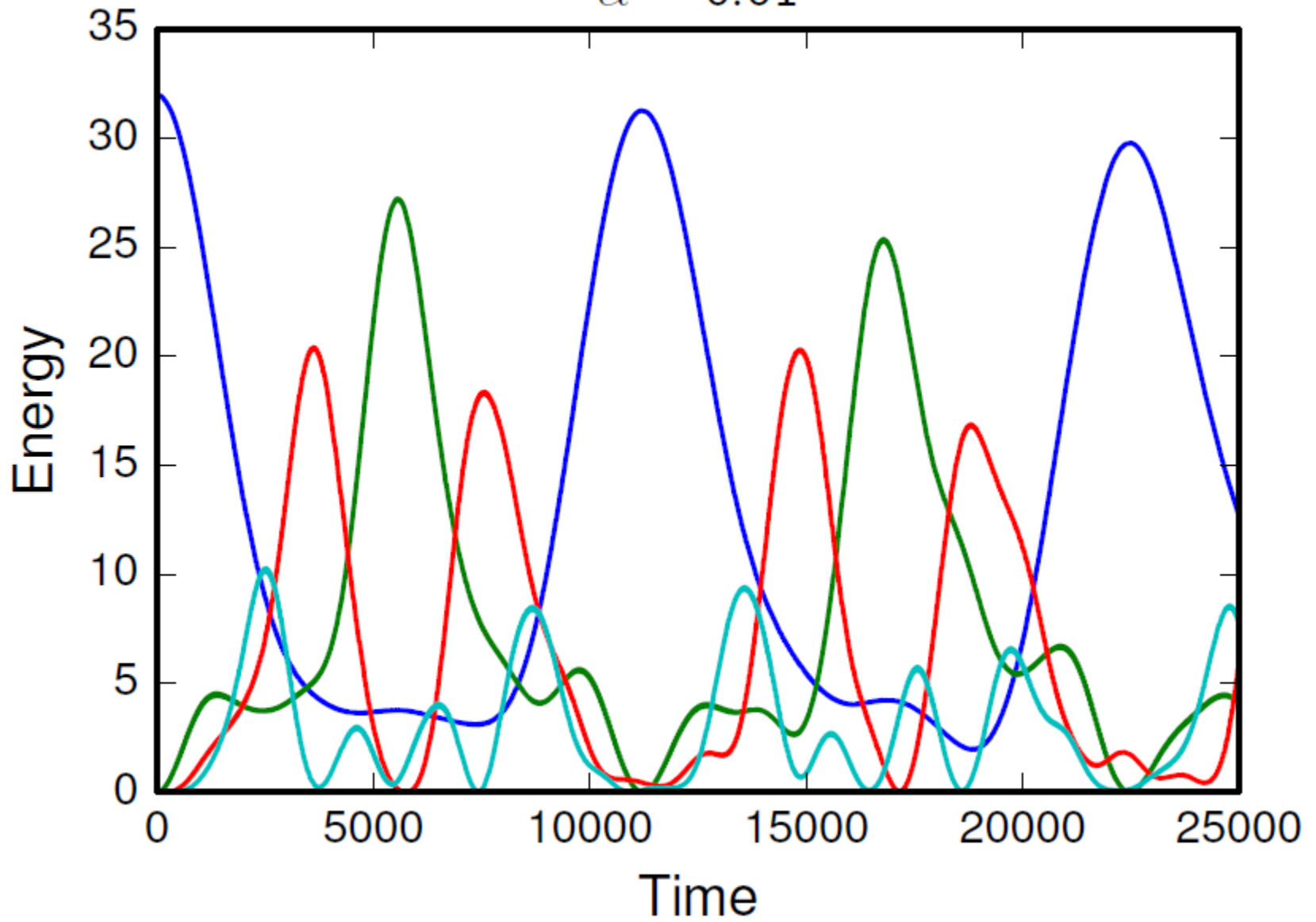
Equipartition of the energy !

Anharmonic system – numerical result



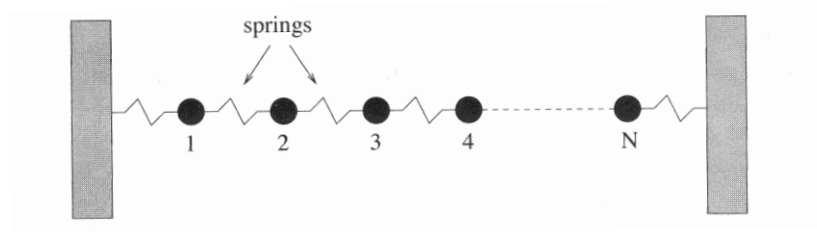
No equipartition of the energy !

$\alpha = 0.01$



Fermi-Pasta-Ulam model

One-dimensional chain of masses connected with non-linear springs.



α -model

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Report written by Fermi, Pasta, and Ulam
Work done by Fermi, Pasta, Ulam, and Tsingou