	Lectures	Exercises/Home work problems	
Ordinary differential equations	Linear dynamics	E1	
	Non-linear dynamics	E2	
	Molecular dynamics		H1a/H1b
Stochastic methods	Monte Carlo integration	E3	
	Metropolis algorithm		H2a/H2b
	Brownian dynamics	E4	
Partial differential equations	Quantum structure	E5	H3a
	Quantum dynamics		H3b

Ising model

Content:

- Ising model
- Ising model: Mean field solution
- Ising model: Metropolis algorithm
- Ising model: Numerical results
- Binary alloy: CuZn
- Binary alloy: Mean field solution

Ising model

The Hamiltonian

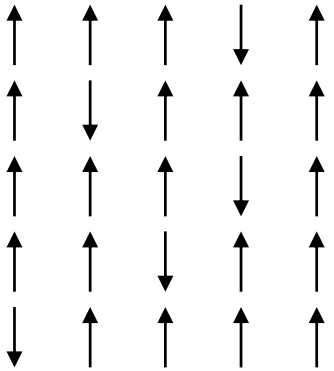
$$\mathcal{H} = -J \sum_{<ij>} s_i s_j - h \sum_i s_i$$

$s_i = \pm 1, \ i = 1, \ldots, N$

$< ij >$ - sum of all nearest neighboring pair of spins

J - coupling constant

h - external field

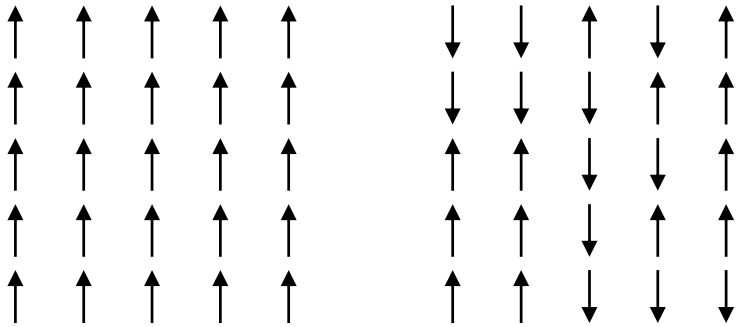


- a lattice model
- one of the simpliest and non-trivial model systems of interacting degrees of freedom
- introduced by Lenz and Ising to model phase transitions in magnetic materials in the 1920s
- solved exactly in 2D by Onsager 1944
- has not been solved exactly in 3D (yet)
- useful in condensed matter physics and field theory
- a (rather crude) model for magnetism
- can be used to model binary alloys in materials science
- can be used to model adsorbed particles in surface science
- can be extended in many directions

Ferromagnetic substance

low T

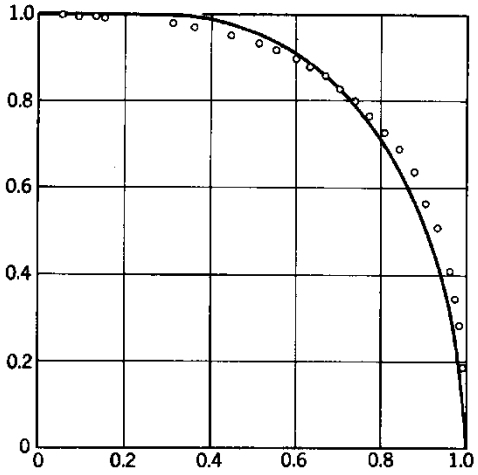
high T



ferromagnetic

paramagnetic

Magnetization

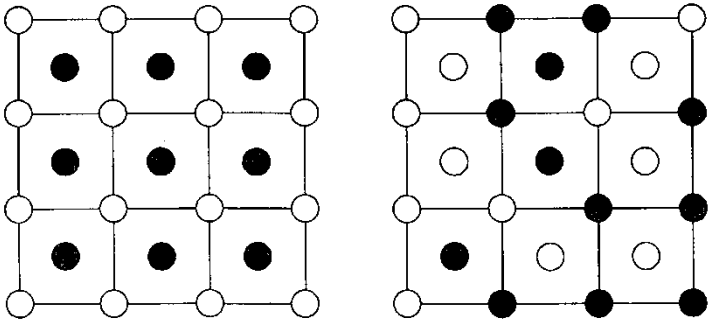


temperature

Binary alloy

low T

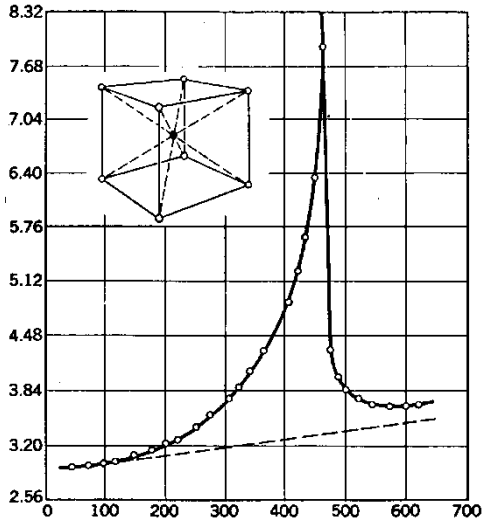
high T



ordered state

disordered state

Heat capacity



temperature

Ising model – statistical thermodynamics

Canonical ensemble:

The probability for the system to be in microstate ν:

$$P_\nu = \frac{1}{Z}\exp(-\beta E_\nu)$$

where $\beta = 1/k_B T$.

The partition function:

$$Z(\beta, h) = \sum_\nu \exp(-\beta E_\nu)$$

The magnetization

$$M_\nu = \sum_{i=1}^{N} s_i$$

The energy

$$E_\nu = -J \sum_{<ij>} s_i s_j - h \sum_i s_i$$

The mean magnetization

$$M \equiv \langle M_\nu \rangle = \frac{1}{Z}\sum_\nu M_\nu \exp(-\beta E_\nu)$$

The mean energy

$$U \equiv \langle E_\nu \rangle = \frac{1}{Z}\sum_\nu E_\nu \exp(-\beta E_\nu)$$

The isothermal susceptibility

$$\chi_T \equiv \left(\frac{\partial M}{\partial h}\right)_T = \frac{1}{k_B T}\left[\left\langle M_\nu^2 \right\rangle - \langle M_\nu \rangle^2\right]$$

The heat capacity at constant field

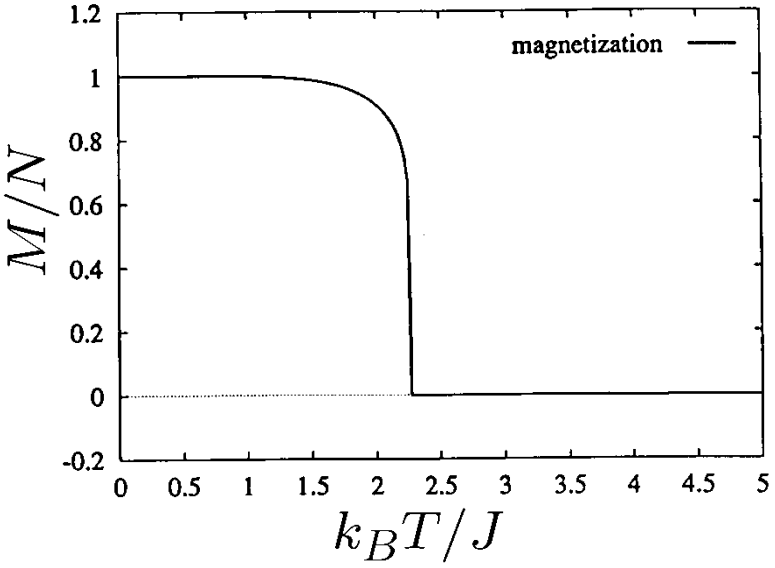
$$C_h \equiv \left(\frac{\partial U}{\partial T}\right)_h = \frac{1}{k_B T^2}\left[\left\langle E_\nu^2 \right\rangle - \langle E_\nu \rangle^2\right]$$

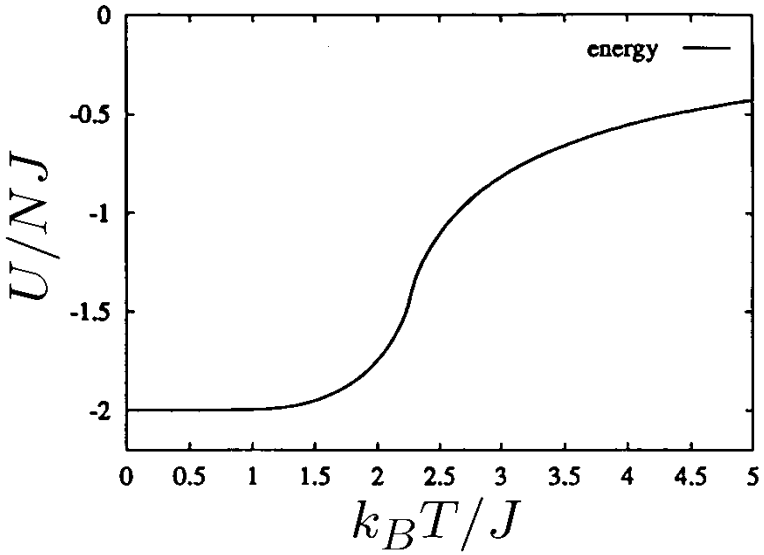
Ising model – the exact solution (2D)

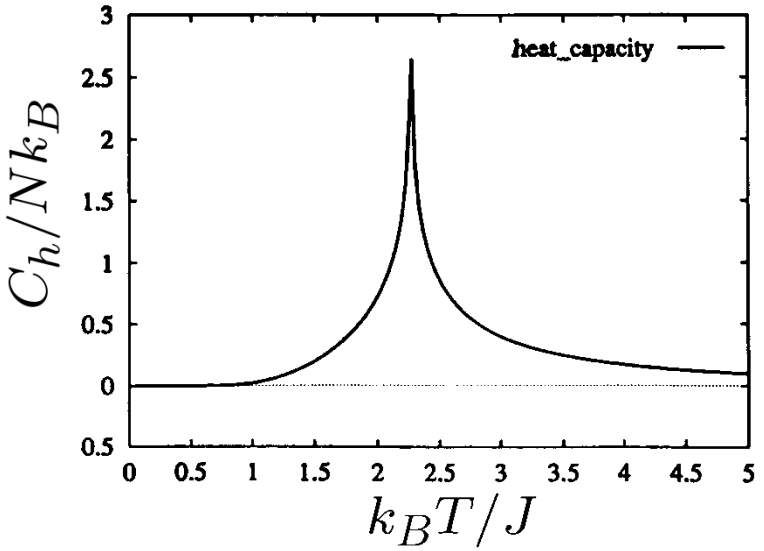
Due to Onsager, 1944

The transition temperature T_c

$$\frac{k_B T_c}{J} = \frac{2}{\ln(\sqrt{2}+1)} \approx 2.269$$







Ising model – mean field solution (h=0)

The exact solution

The energy

$$U = -J \sum_{<ij>} \left\langle s_i s_j \right\rangle$$

The mean field solution

The energy

$$U_{MF} = -J \sum_{<ij>} \langle s_i \rangle \langle s_j \rangle = -J 2N m^2$$

The number of microstates

$$W = \frac{N!}{N_\uparrow! N_\downarrow!}$$

The entropy

$$S = k_B \ln W = \frac{N!}{N_\uparrow! N_\downarrow!}$$

The free energy

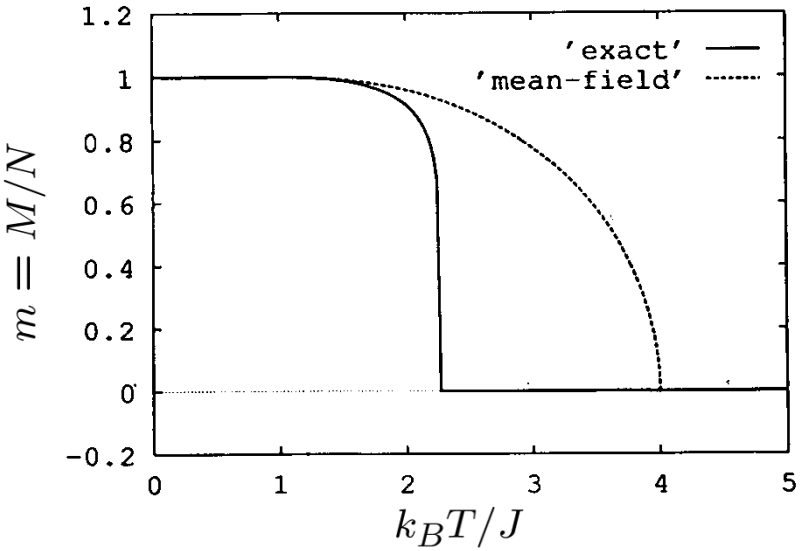
$$F = U - TS$$

Minimum if

$$m = \frac{k_B T}{8J} \ln \frac{1+m}{1-m}$$

which implies that

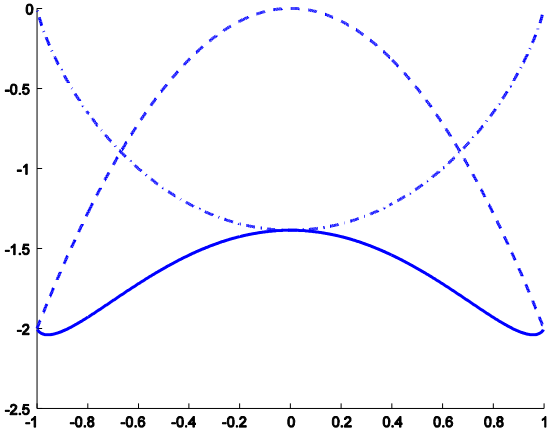
$$k_B T_c / J = 4.0$$

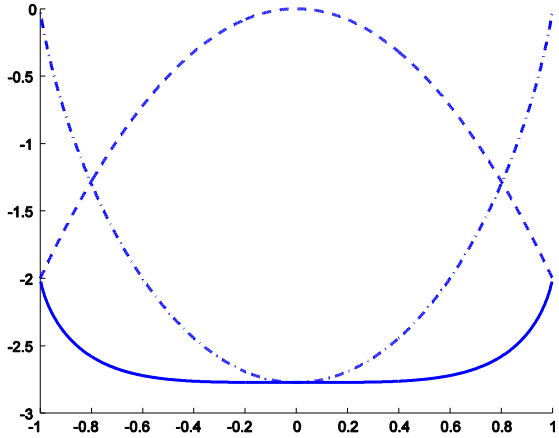


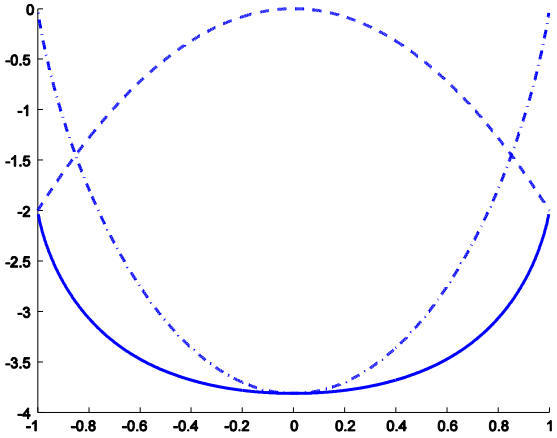
Ising model – mean field solution (h=0)

$$\frac{F_{MF}}{NJ} = -2m^2 + \frac{k_B T}{J}\left[\frac{1+m}{2}\ln\frac{1+m}{2} + \frac{1-m}{2}\ln\frac{1-m}{2}\right]$$

$k_B T/J = 2.5$ $k_B T/J = k_B T_c/J = 4.0$ $k_B T/J = 5.5$







Dashed line: U

Dash-dotted line: $-TS$

Full line: $F = U - TS$

Ising model – Metropolis algorithm

1. Set the desired temperature T and external field h.

2. Initialize the system, *e.g.* use a random configuration or a configuration from a previous simulation.

3. Perform the desired number of Monte Carlo sweeps through the lattice.

4. Exclude the first configurations (let the system equilibrate).

5. Compute average quantities from subsequent configurations and estimate the error from statistically independent configurations.

3a. Make a trial change, e.g. by flipping a randomly chosen spin.

3b. Determine the change in energy ΔE

3c. If $\Delta E \leq 0$ accept the new configuration

3d. If $\Delta E > 0$ generate a random number r between 0 and 1, and if

$$\exp(-\Delta E/k_B T) \geq r$$

accept the new configuration, otherwise count the old configuration once more.

The Metropolis algorithm

The Metropolis algorithm is a particular way of ensuring that the transition rule satisfies detailed balance.

The transition matrix $w_{nm} = w_{n \leftarrow m}$ is split into two parts

$$w_{nm} = \tau_{nm} \alpha_{nm}$$

where τ_{nm} is the probability of making a trial change from state Ω_m to state Ω_n and α_{nm} is the probability of accepting the trial state.

The acceptance probability is assumed to satisfy

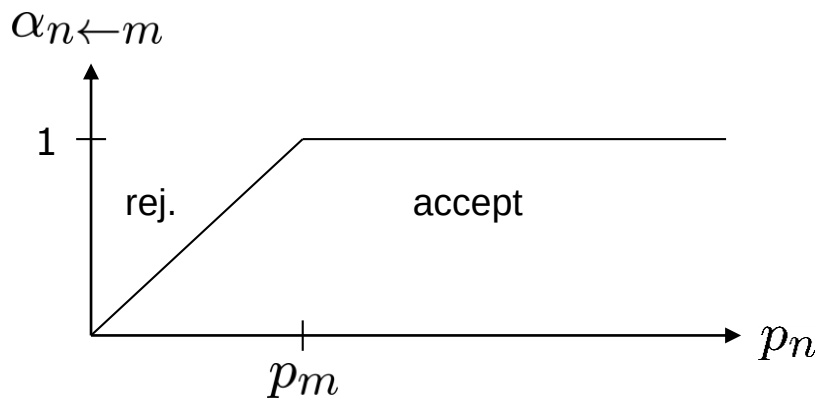
$$\alpha_{nm} = \begin{cases} 1 & \text{if } p_n \geq p_m \\ p_n / p_m & \text{if } p_n < p_m \end{cases}$$

and to ensure detailed balance the trail change then has to be symmetric

$$\tau_{nm} = \tau_{mn}$$

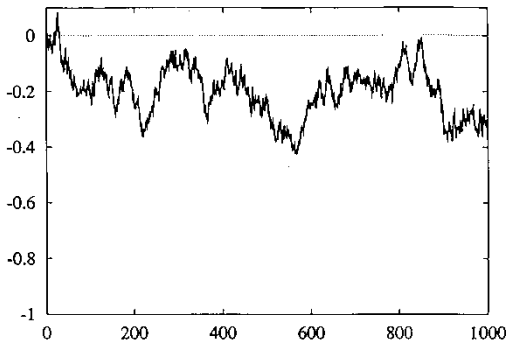
Moves that are not accepted are rejected and remain at the same location for at least one more step

$$w_{mm} = 1 - \sum_{n(\neq m)} w_{nm}$$



Equilibration

It is important to wait until
the system has **equilibrated.**



Equilibration

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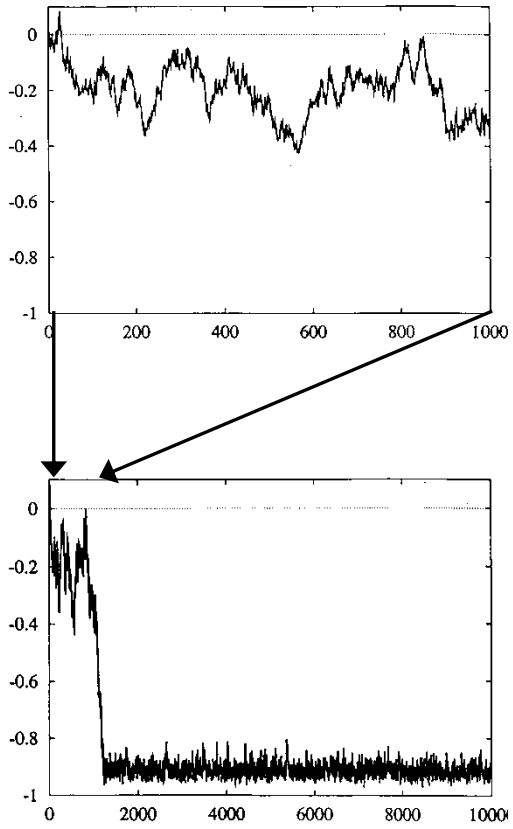

Equilibration

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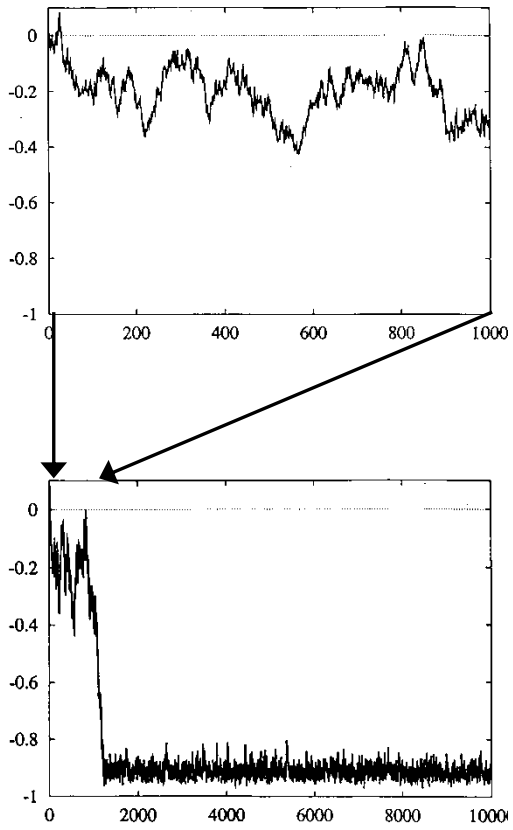
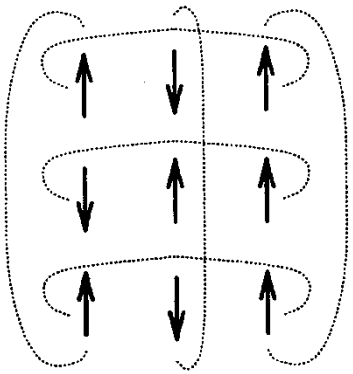
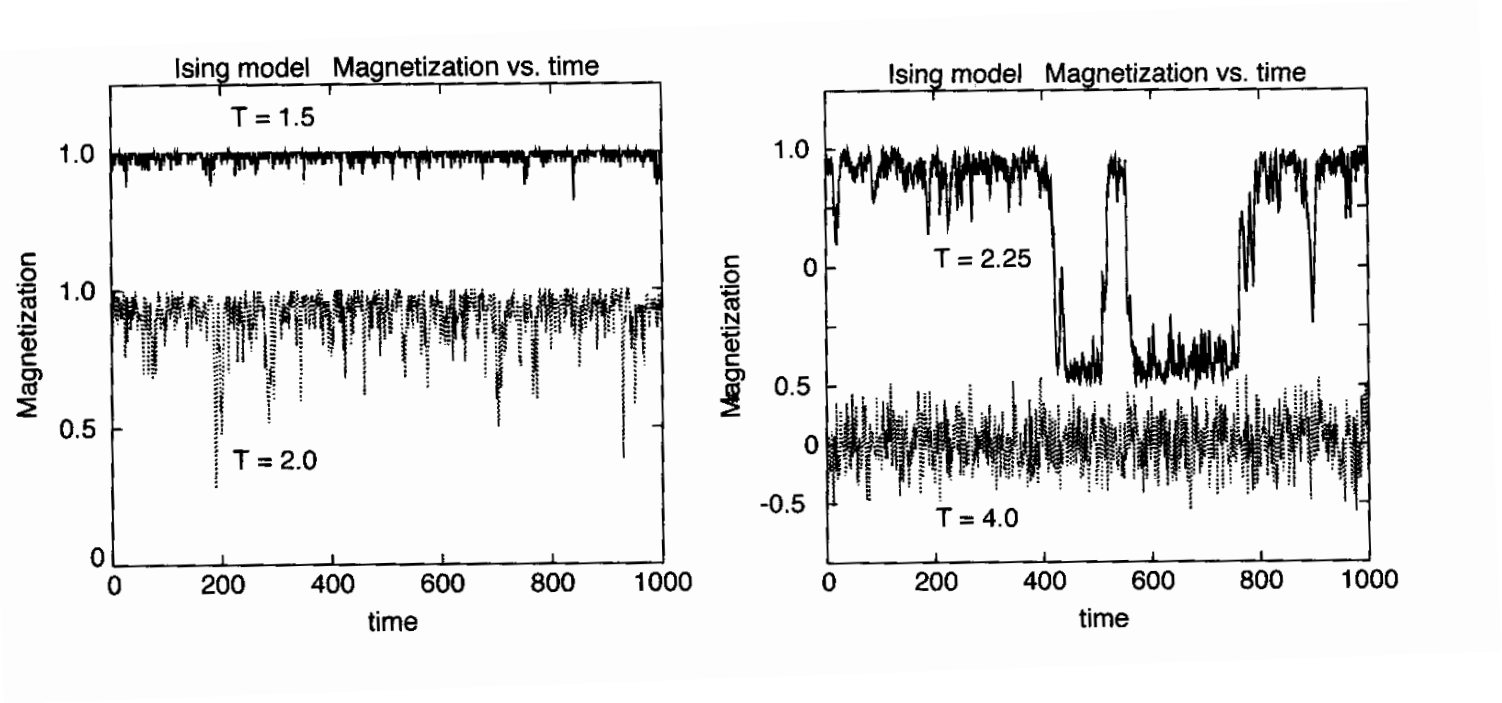


Boundary conditions

To mimic a large system **periodic boundary conditions** are commonly used.

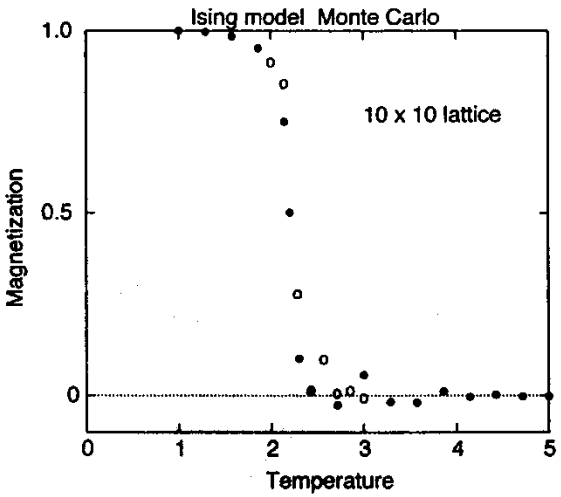


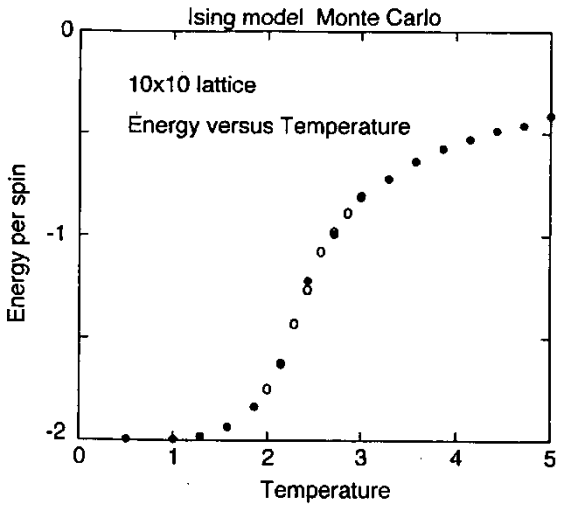
Ising model – Numerical results

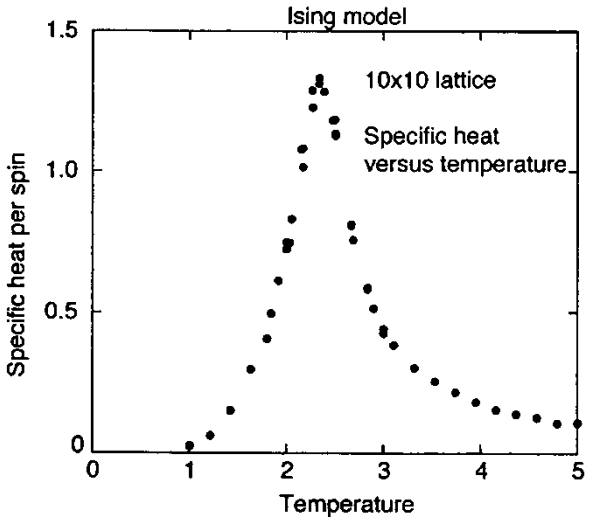


The temperature T in units of J/k_B.

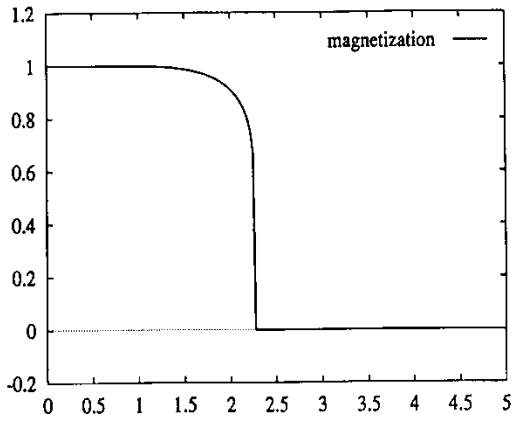
Ising model – Numerical results

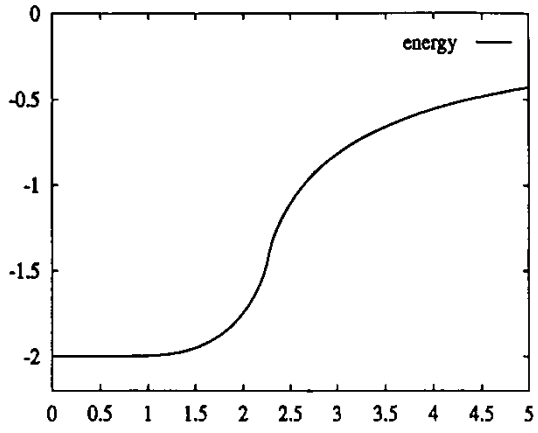


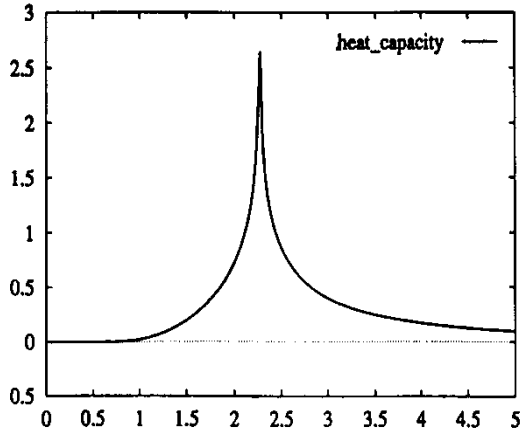




Ising model – Exact results



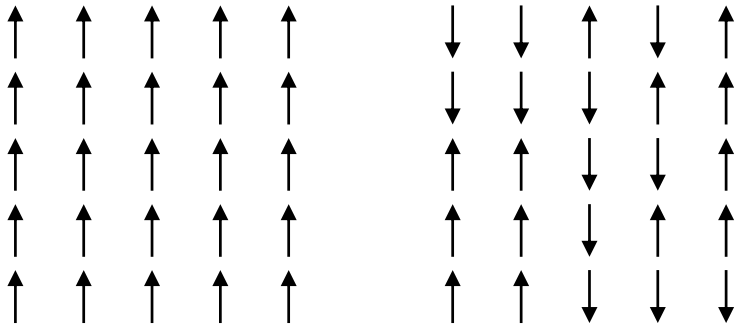




Ferromagnetic substance

low T

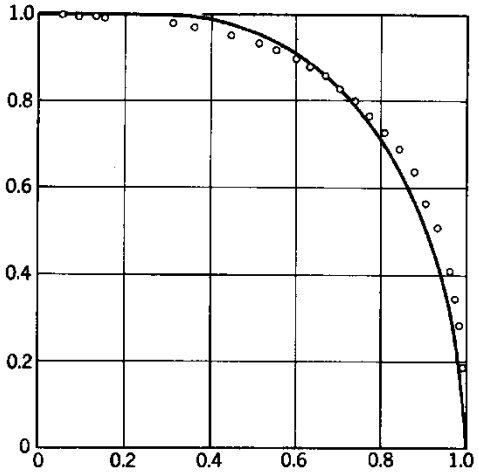
high T



ferromagnetic

paramagnetic

Magnetization

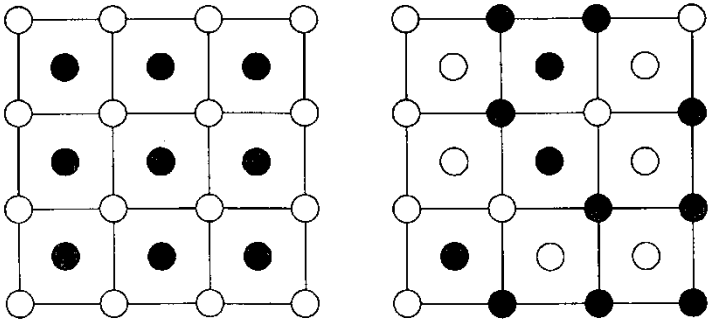


temperature

Binary alloy

low T

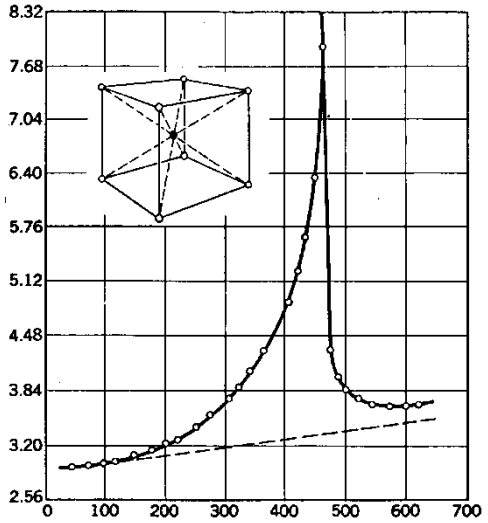
high T



ordered state

disordered state

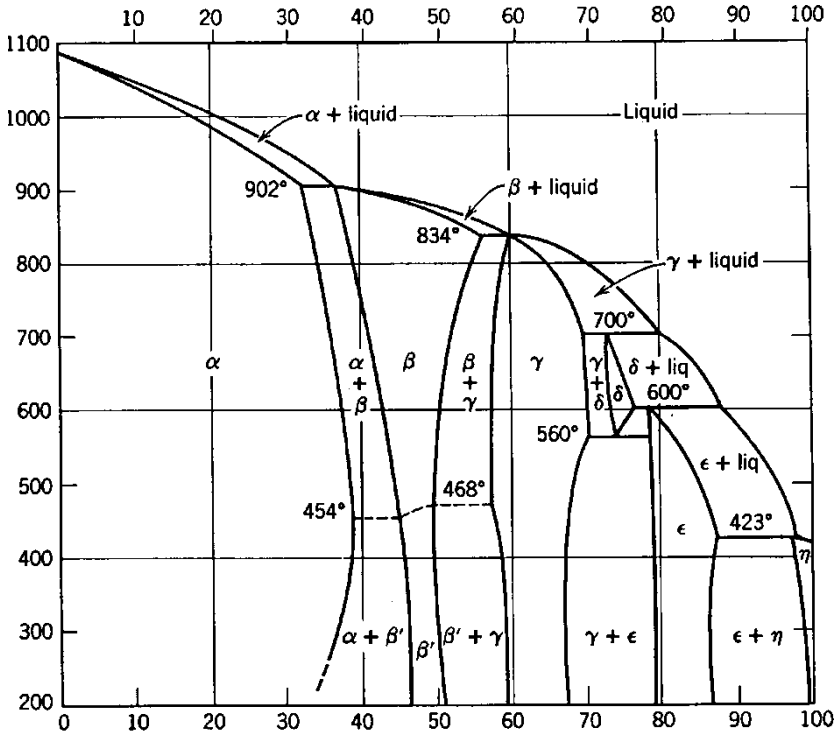
Heat capacity



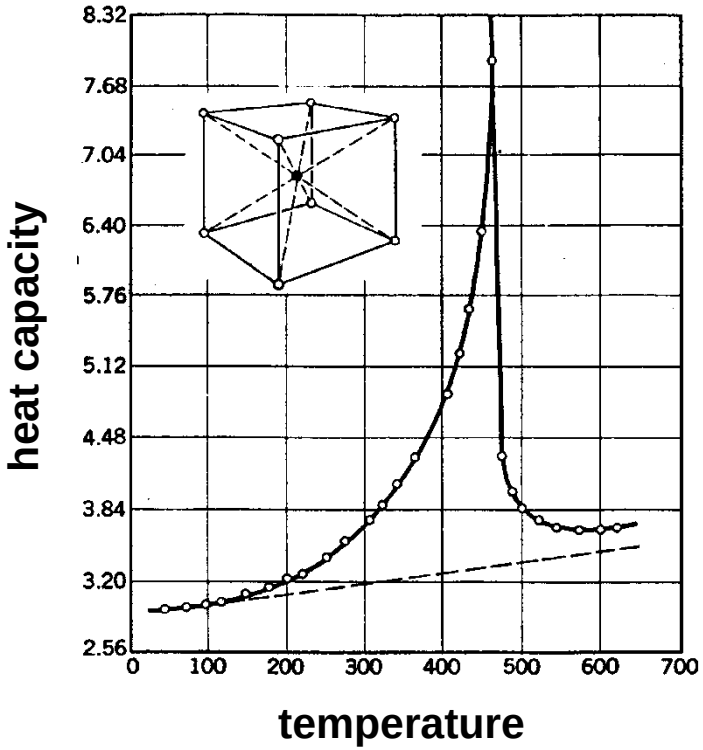
temperature

The Cu-Zn system – a binary alloy

Equilibrium phase diagram

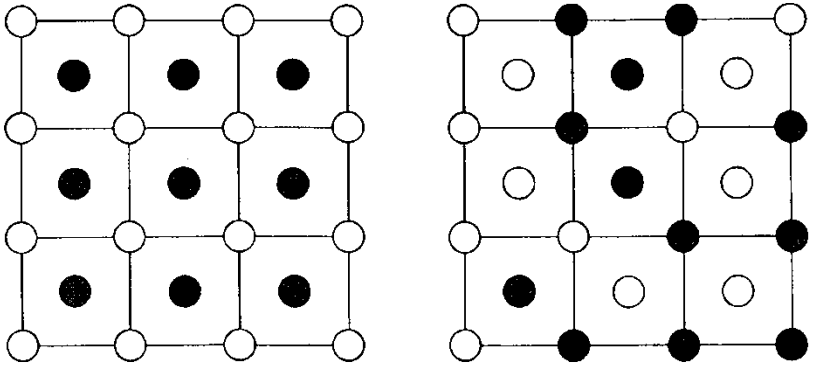


CuZn (β-brass)



Binary alloy – simple model

AB alloy with bcc structure
ex. Cu-Zn (β-brass)



- N atoms A, N atoms B

- bcc-lattice = two interpenetrating sc-lattices (a and b)

Long-range order parameter P

Number of A atoms on the a sublattice
$= \frac{1}{2}(1+P)N$

$P = \pm 1$ perfect order
$P = 0$ no order

Short-range order parameter r

q = number of nearest-neighbor bonds that are AB bonds

$$r = \frac{1}{4}(q-4)$$

$r = 1$ complete order
$r = 0$ complete disorder

Binary alloy – mean field solution

The energy

$$E = N_{AA}E_{AA} + N_{BB}E_{BB} + N_{AB}E_{AB}$$

where N_{ij} is the number of nearest-neighbor ij bonds, and E_{ij} is the energy of an ij bond.

Mean field approximation

Assume no correlations, i.e.

$$\begin{aligned} N_{AA} &= 8\,[\tfrac{1}{2}(1+P)N]\,[\tfrac{1}{2}(1-P)] = 2(1-P^2)N \\ N_{BB} &= 8\,[\tfrac{1}{2}(1+P)N]\,[\tfrac{1}{2}(1-P)] = 2(1-P^2)N \\ N_{AB} &= 8N\,[\tfrac{1}{2}(1+P)]^2 + 8N\,[\tfrac{1}{2}(1-P)]^2 = 4(1+P^2)N \end{aligned}$$

This implies that

$$E_{MF} = E_0 - 2NP^2\Delta E$$

where

$$\begin{aligned} E_0 &= 2N(E_{AA} + E_{BB} + 2E_{AB}) \\ \Delta E &= E_{AA} + E_{BB} - 2E_{AB} \end{aligned}$$

Binary alloy – mean field solution

The number of configurations

$$W = \left[\frac{N!}{[\frac{1}{2}(1+P)N]!\ [\frac{1}{2}(1-P)N]!} \right]^2$$

The entropy

$$\begin{aligned} S &= k_B \ln W \\ &= 2Nk_B \ln 2 - Nk_B[(1+P)\ln(1+P) + (1-P)\ln(1-P)] \end{aligned}$$

The free energy

$$\begin{aligned} F &= U - TS \\ &= E_0 - 2NP^2\Delta E \\ &- 2Nk_BT \ln 2 + Nk_BT[(1+P)\ln(1+P) + (1-P)\ln(1-P)] \end{aligned}$$

The equilibrium structure is obtained by finding the minimum of F with respect to the order parameter P. It leads to a phase transition at

$$T_c = 2\Delta E/k_B$$

References:

Computational treatment of the Ising model:

- Giordano and Nakanishi, *Computational Physics*
- Koonin and Meredith, *Computational Physics*

Binary alloy:

- Kittel, *Introduction to Solid State Physics*

Advanced treatment of the Ising model:

- Huang, *Statistical Mechanics*

Variational Monte Carlo

Content:

- Variational Quantum Monte Carlo
- Multi-dimensional case
- Example: Helium

Variational Quantum Monte Carlo

Aim: To obtain a good estimate of the ground state energy E_0 for a quantum system.

Based on the variational theorem

$$E[\psi_T] = \frac{\langle \psi_T \mid \mathcal{H} \mid \psi_T \rangle}{\langle \psi_T \mid \psi_T \rangle} \geq E_0$$

where

ψ_T is a trial wave function

Choose a physically reasonable form for the trial wave wave function that depends on one or more parameters α_i. Determine $E[\psi_T]$ and vary α_i until a minimum is obtained.

Multi-dimensional case

Write in terms of a normalized weight function $\rho(\boldsymbol{X})$, or probability distribution

$$E[\psi_T] = \frac{\int d\boldsymbol{X}\psi_T^*(\boldsymbol{X})\mathcal{H}\psi_T(\boldsymbol{X})}{\int d\boldsymbol{X}\psi_T^*(\boldsymbol{X})\psi_T(\boldsymbol{X})} = \int d\boldsymbol{X}E_L(\boldsymbol{X})\rho(\boldsymbol{X})$$

where

$$E_L(\boldsymbol{X}) = \frac{\mathcal{H}\psi_T(\boldsymbol{X})}{\psi_T(\boldsymbol{X})}$$

is the called the local energy and

$$\rho(\boldsymbol{X}) = \frac{\mid \psi_T(\boldsymbol{X}) \mid^2}{\int d\boldsymbol{X} \mid \psi_T(\boldsymbol{X}) \mid^2}$$

is the weight function. Use the Metropolis algorithm to evaluate the multi-dimensional integral.

Example: Helium

Hamiltonian in atomic units ($\hbar = m_e = e = 4\pi\epsilon_0 = 1$)

$$\mathcal{H} = -\frac{1}{2}(\nabla_1^2 + \nabla_2^2) - \frac{2}{r_1} - \frac{2}{r_2} + \frac{1}{r_{12}}$$

Trial wavefunction

$$\psi_T(\boldsymbol{r}_1, \boldsymbol{r}_2) = \phi(r_1)\phi(r_2)f(r_{12})$$

with

$$\begin{aligned} \phi(r) &= \exp\left[-\alpha_1 r\right] \\ f(r) &= \exp\left[\frac{\alpha_2 r}{1+\alpha_3 r}\right] \end{aligned}$$

Example: Helium

The cusp conditions (boundary conditions)

$$\lim_{r\to 0}\frac{1}{\phi}\frac{\partial\phi(r)}{\partial r} = -2$$

$$\lim_{r\to 0}\frac{1}{f}\frac{\partial f(r)}{\partial r} = \frac{1}{2}$$

imply that

$$\psi_T(\boldsymbol{r}_1,\boldsymbol{r}_2) = \exp\left[-2r_1\right]\ \exp\left[-2r_2\right]\ \exp\left[\frac{r_{12}}{2(1+\alpha r_{12})}\right]$$

and that the local energy can be written as

$$E_L(\boldsymbol{r}_1,\boldsymbol{r}_2) = -4 + \frac{(\widehat{\boldsymbol{r}}_1-\widehat{\boldsymbol{r}}_2)\cdot\widehat{\boldsymbol{r}}_{12}}{(1+\alpha r_{12})^2}$$
$$+\frac{\alpha}{(1+\alpha r_{12})} + \frac{\alpha}{(1+\alpha r_{12})^2} + \frac{\alpha}{(1+\alpha r_{12})^3} - \frac{1}{4(1+\alpha r_{12})^4}$$

References:

Variational Quantum Monte Carlo:

- Koonin and Meredith, *Computational Physics*
- Thijssen, *Computational Physics*
- P.J. Reynolds, J. Tobochnik, and H. Gould, Computers in Physics, nov/dec 1990.

Green's function and Diffusion Monte Carlo:

- Koonin and Meredith, *Computational Physics*
- Thijssen, *Computational Physics*